

New finding in electronic properties of superconducting cuprates from nodal photoemission spectra

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Experimental collaboration, single crystals

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S. Ono, Yoichi Ando

CRIEPI Tokyo

Band structure calculations

Alexander Yaresko

MPI-PKS Dresden

Stefan-Ludwig Drechsler

IFW Dresden

Navigation

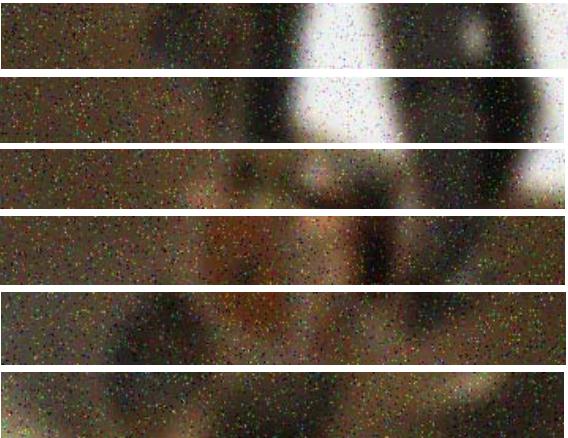
Introduction to ARPES

The advantages of our group

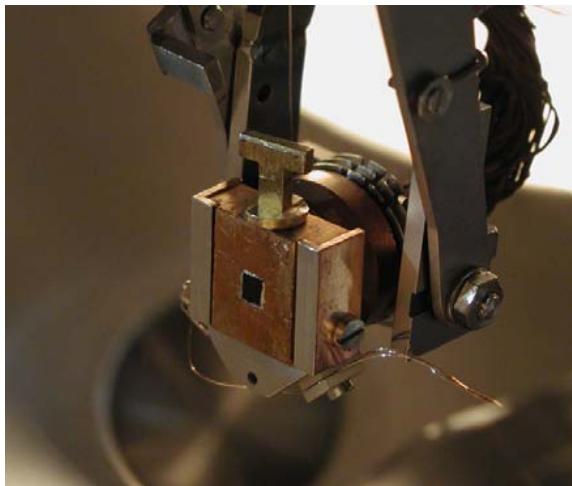
Electronic band structure

Antinodal region

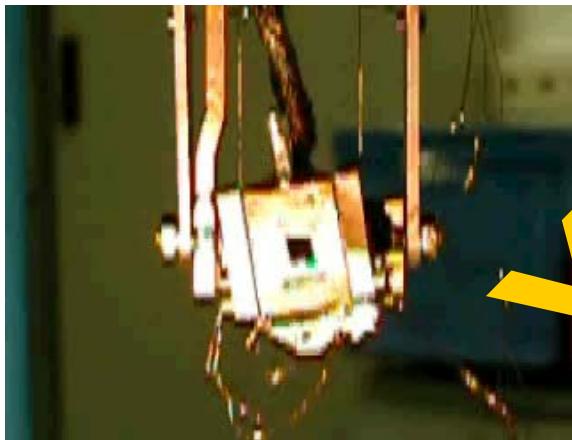
Nodal region



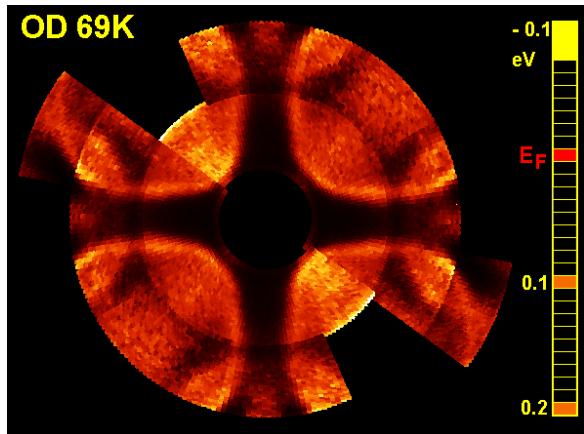
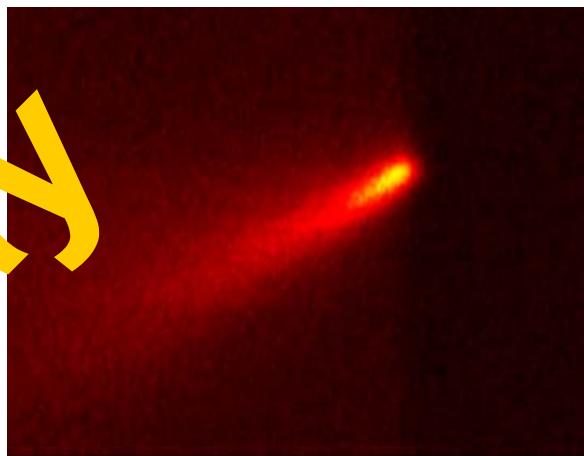
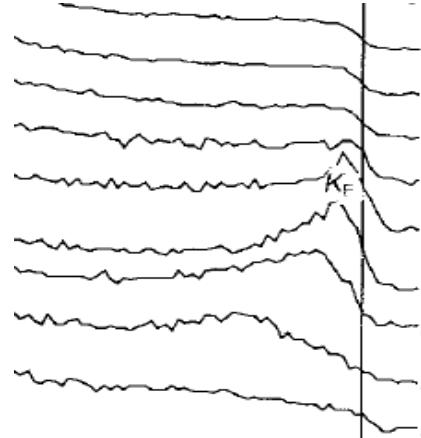
Last century ARPES



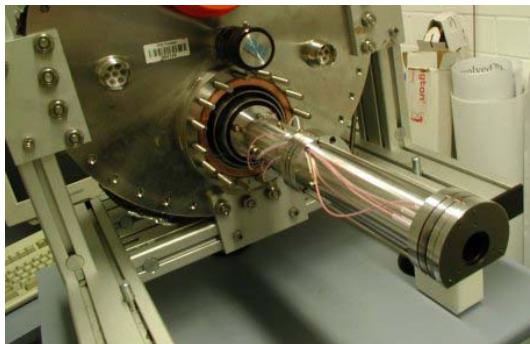
Today



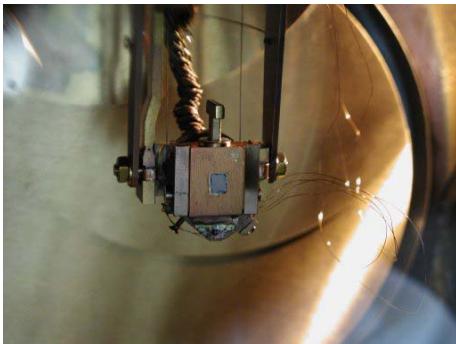
Yesterday
Today in Dresden



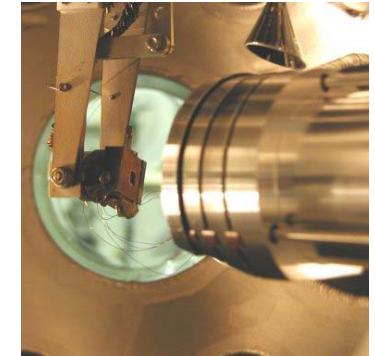
(k,ω) -space explorer today



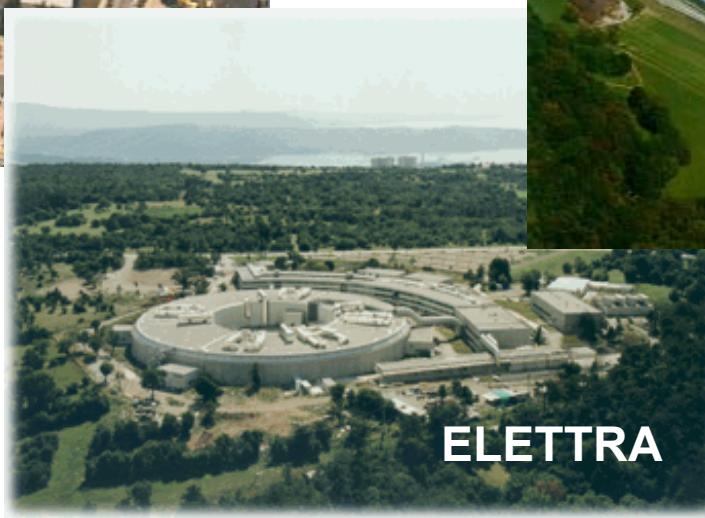
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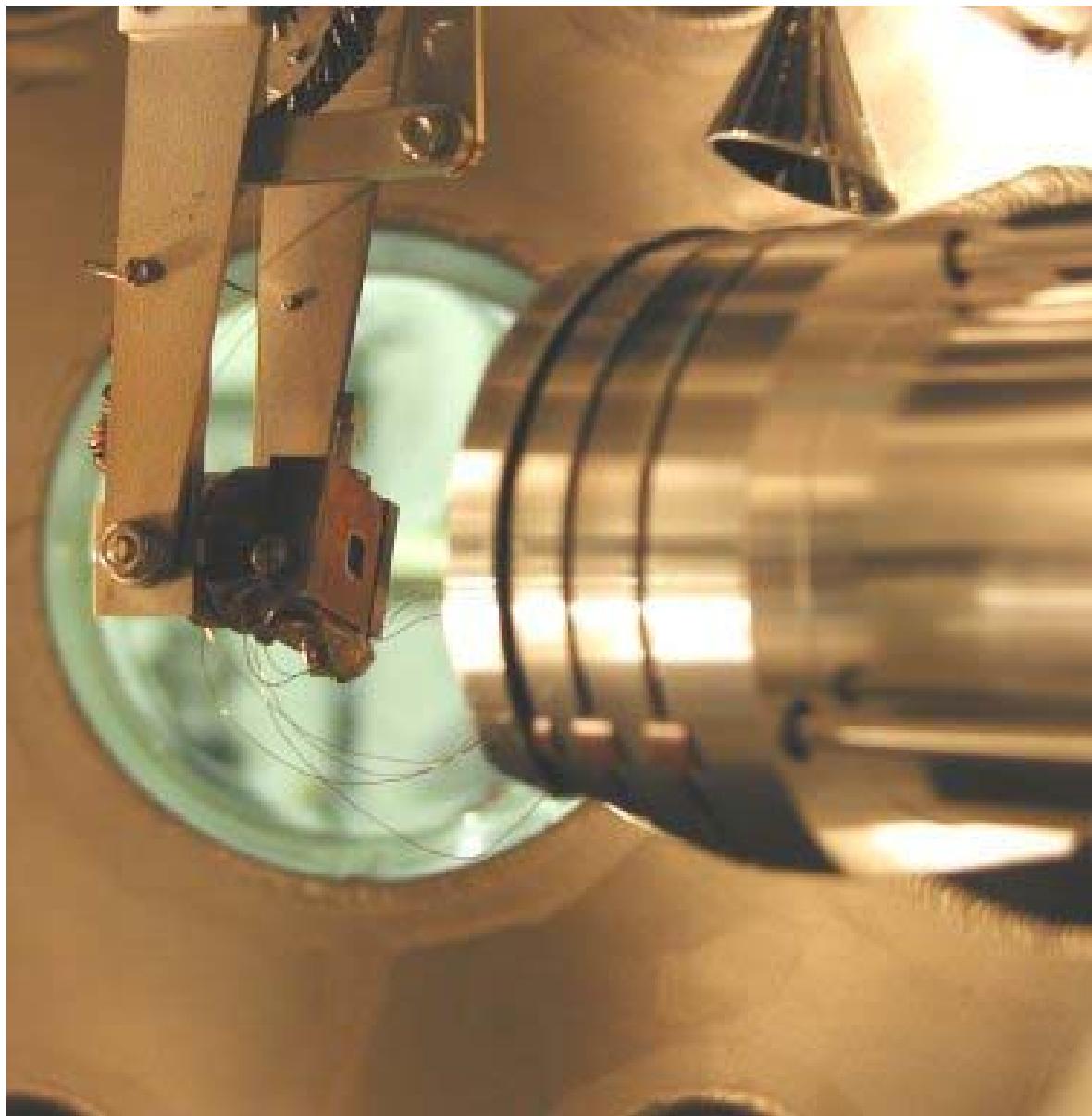
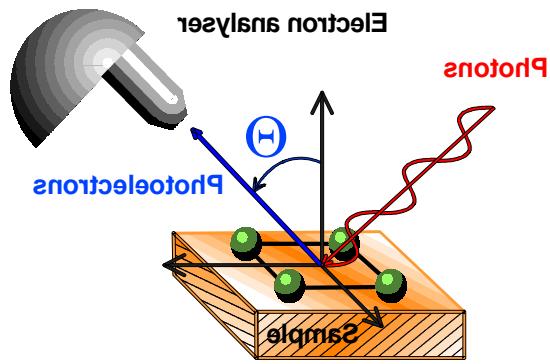
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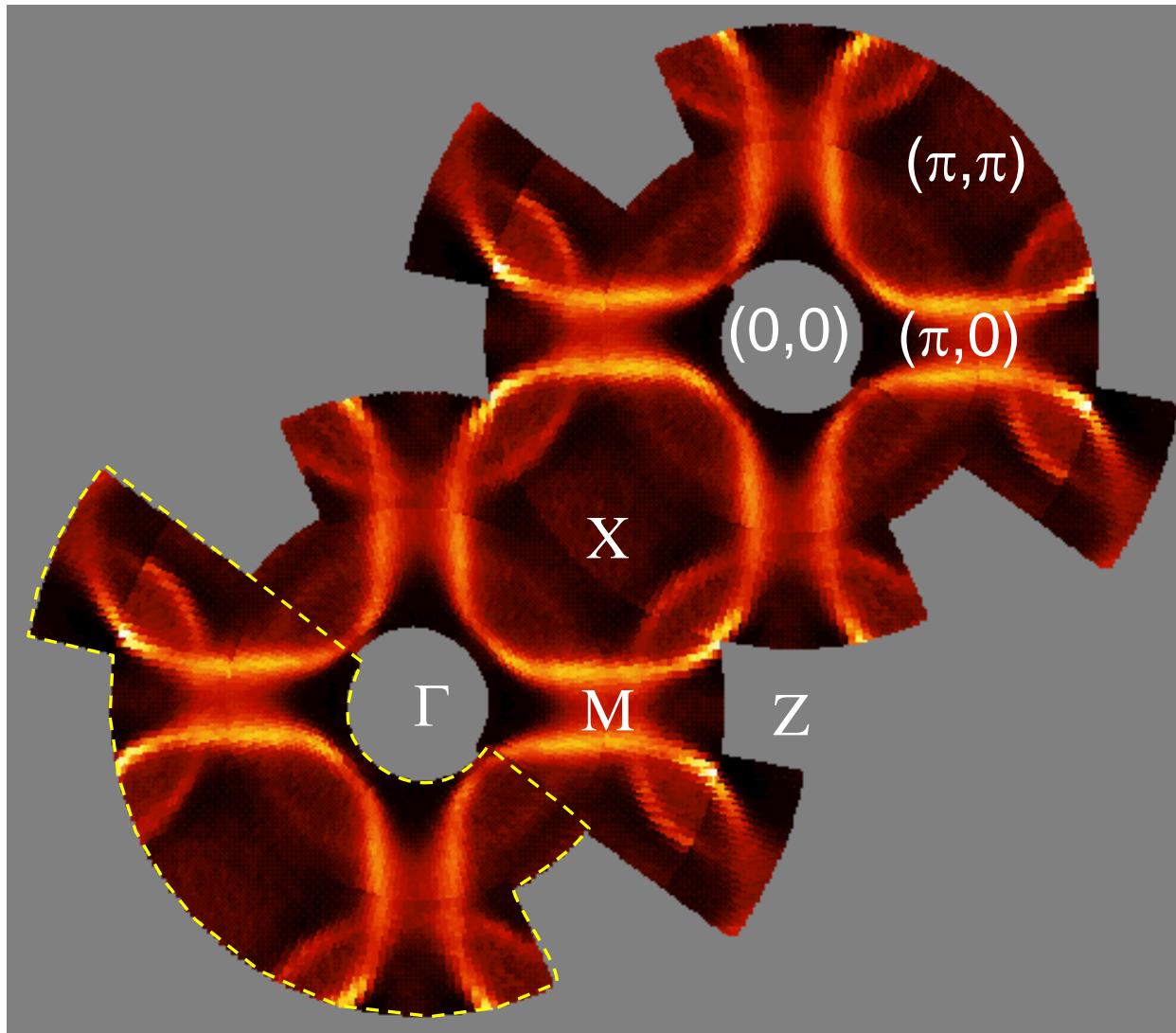
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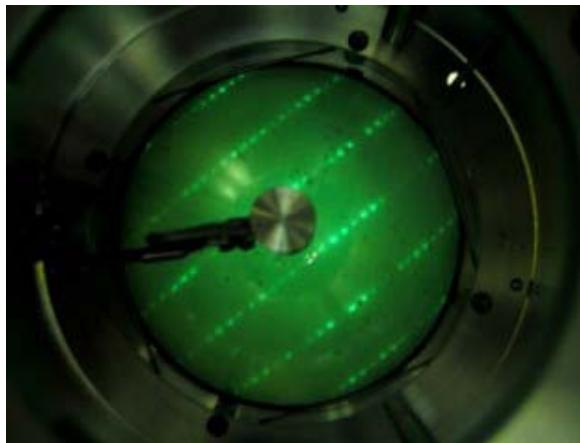
Precise Cryo-Manipulator



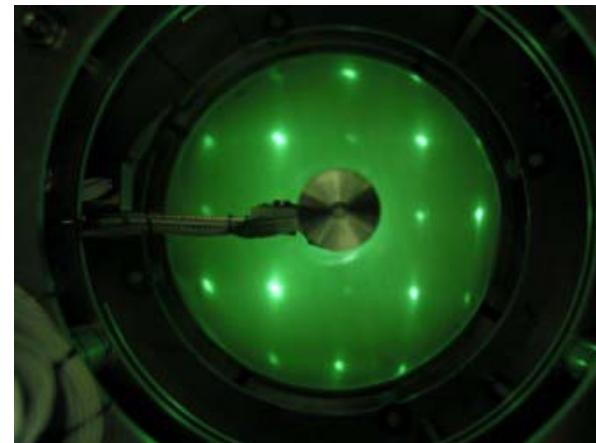
Fermi-surface map



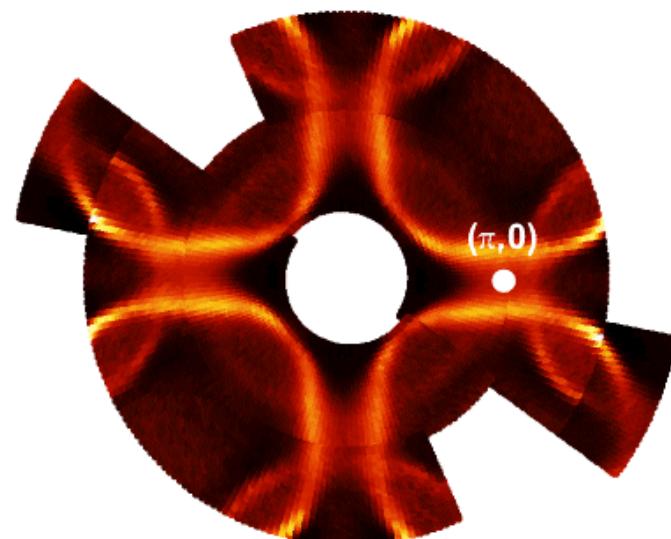
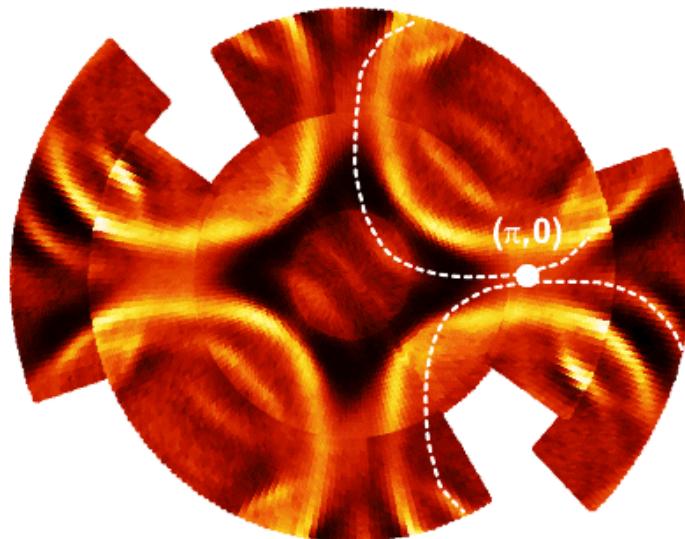
Superstructure free samples...



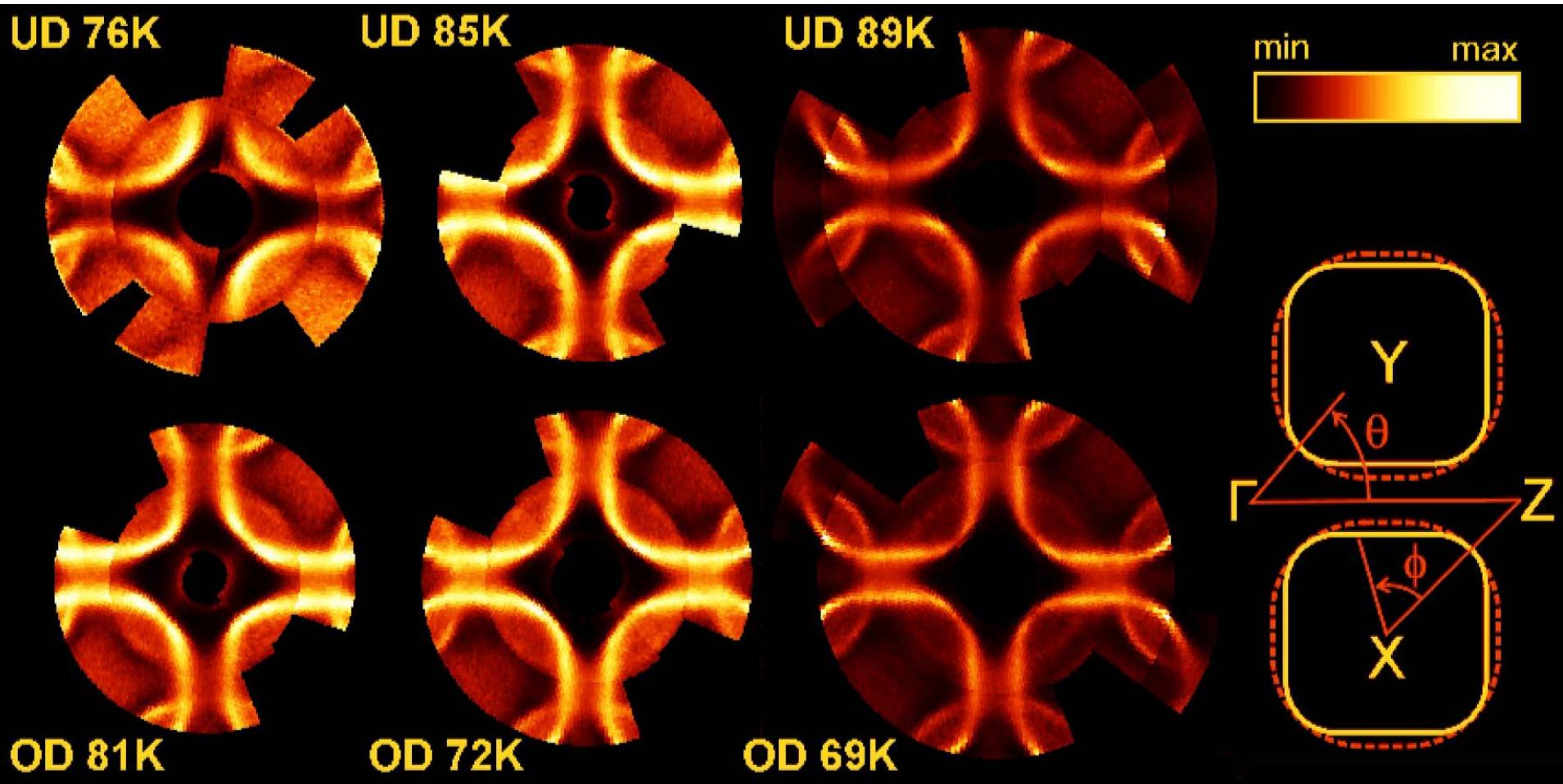
Bi2212



Pb-Bi2212

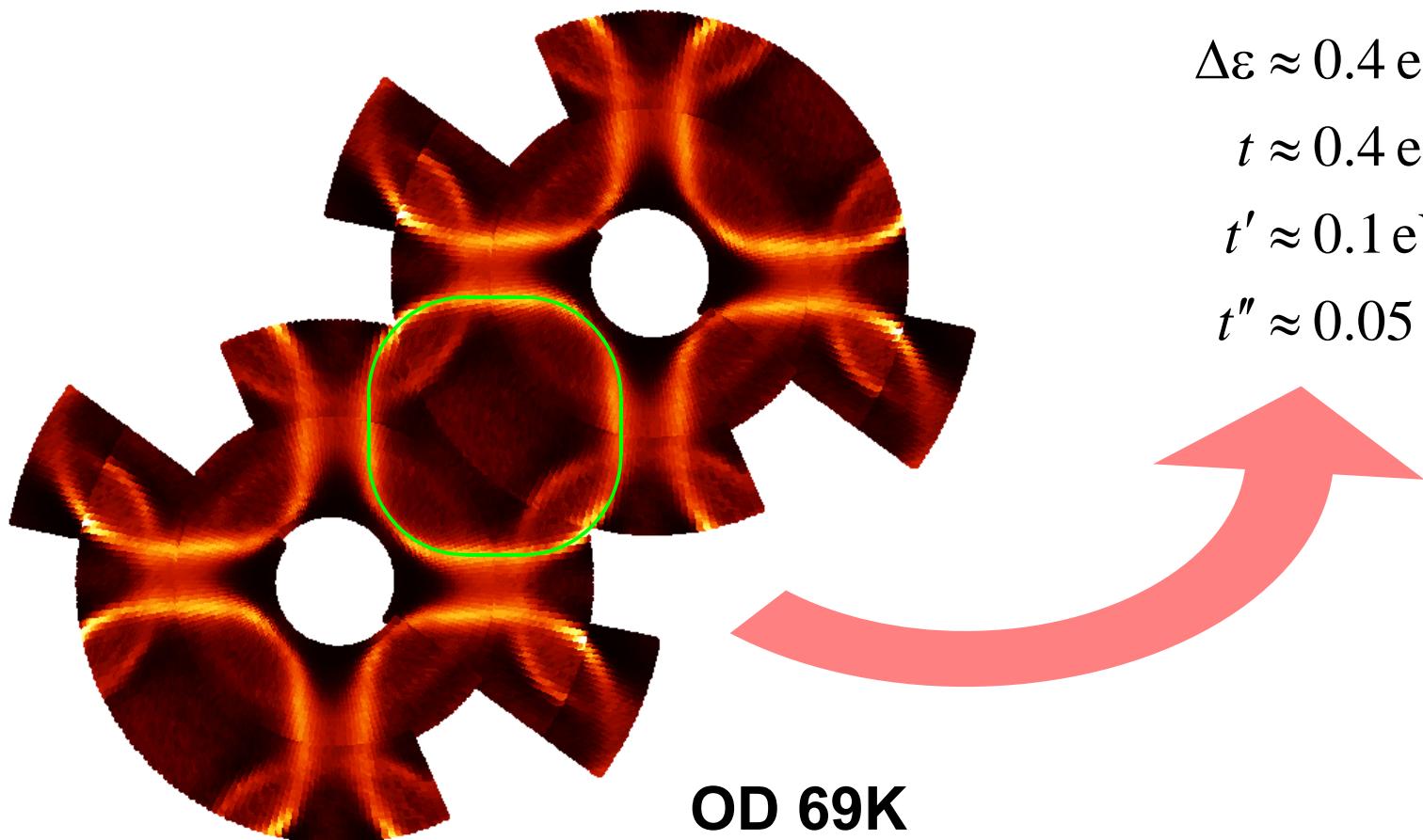


...in a wide doping range



Band structure: TBF

$$\varepsilon(k_x, k_y) = \Delta\varepsilon - 2t(\cos k_x + \cos k_y) + 4t' \cos k_x \cos k_y - 2t''(\cos 2k_x + \cos 2k_y)$$



$$\Delta\varepsilon \approx 0.4 \text{ eV}$$

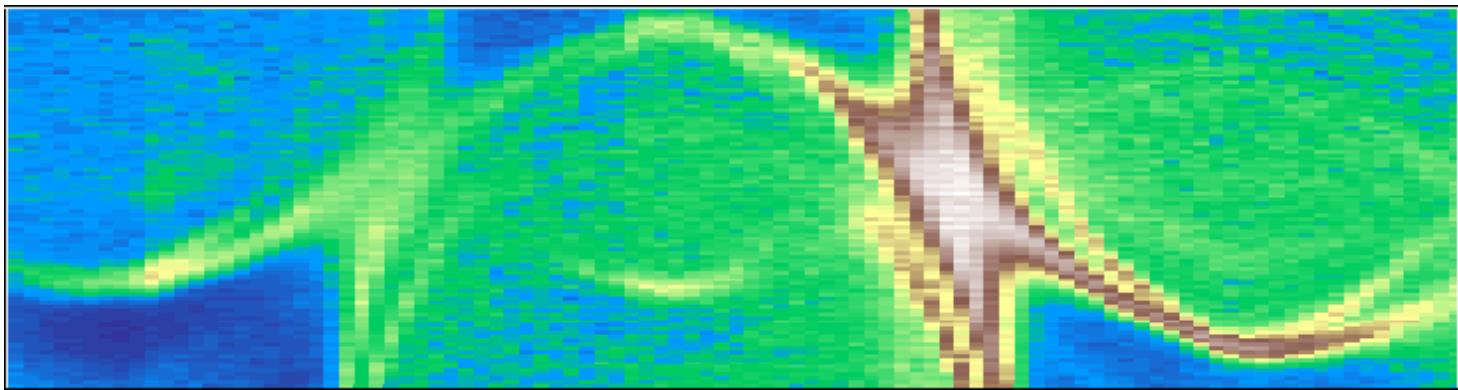
$$t \approx 0.4 \text{ eV}$$

$$t' \approx 0.1 \text{ eV}$$

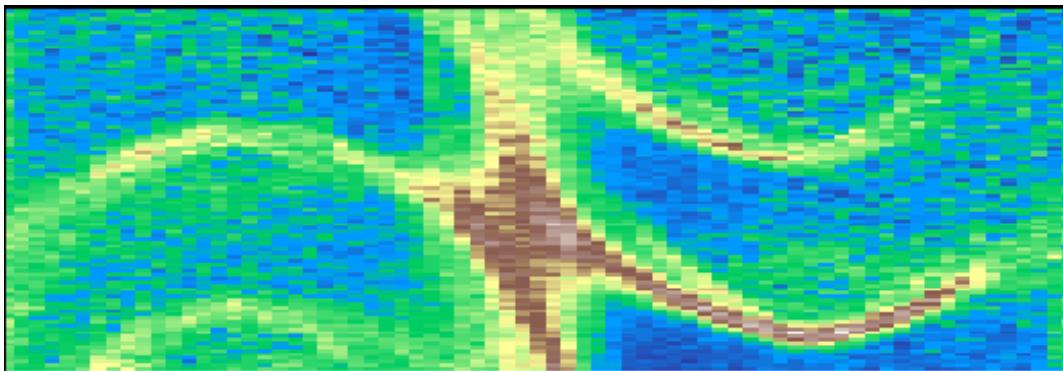
$$t'' \approx 0.05 \text{ eV}$$

High precision Fermi surface mapping

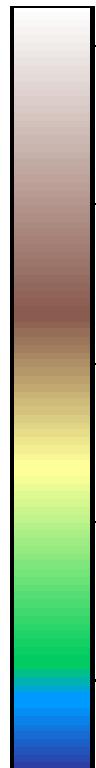
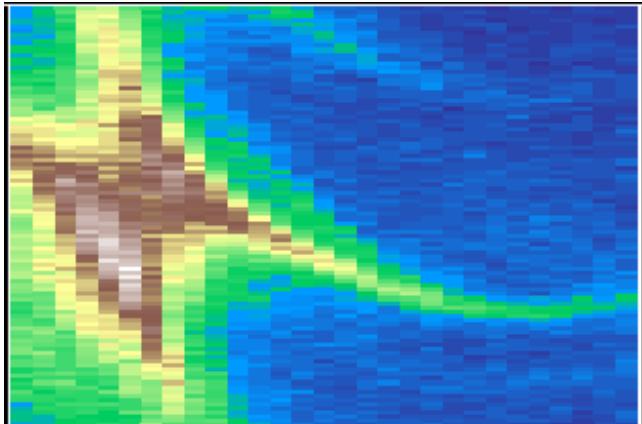
OD
 $< T_c$



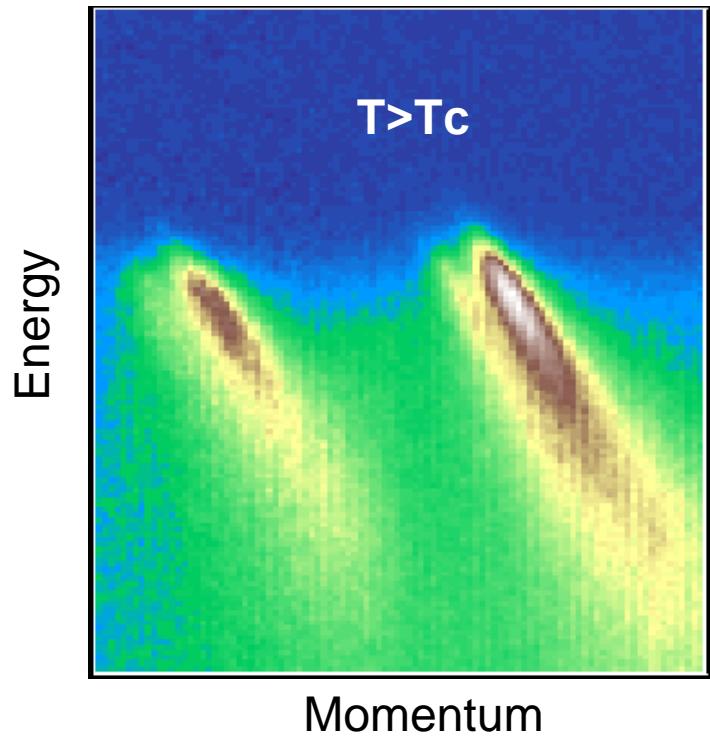
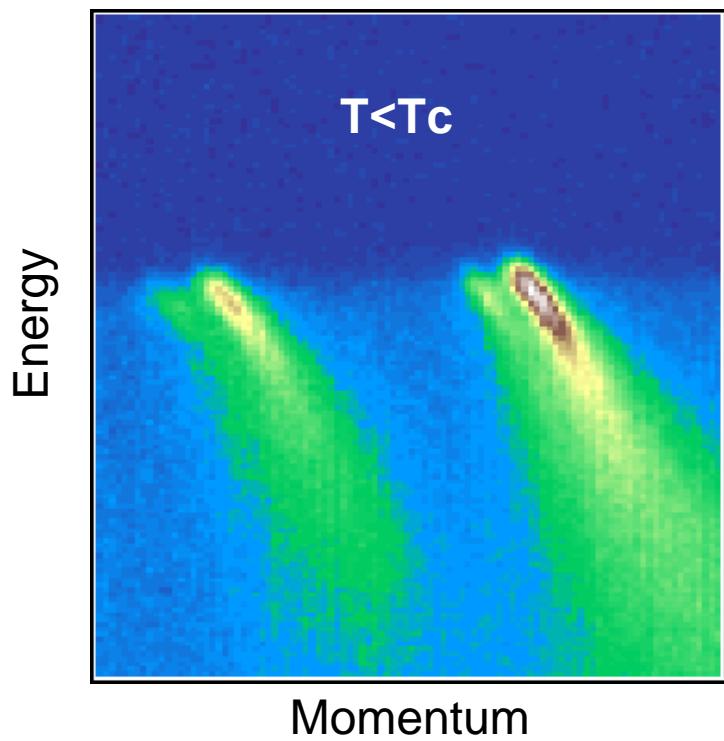
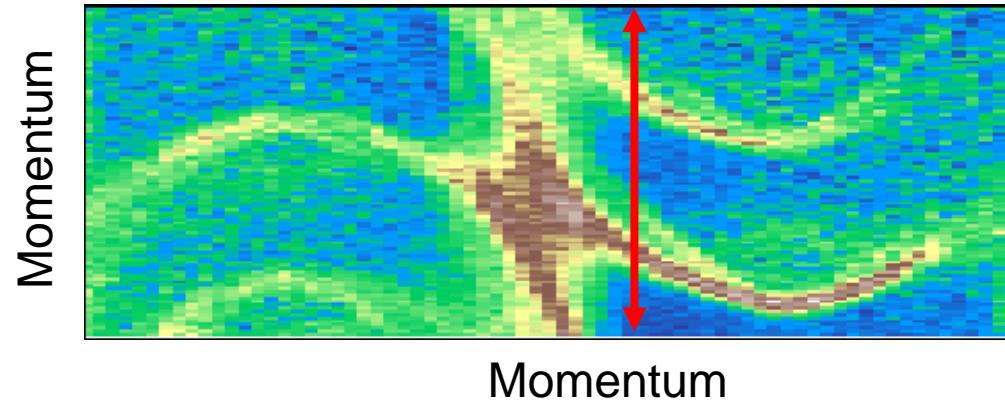
OP
 $< T_c$



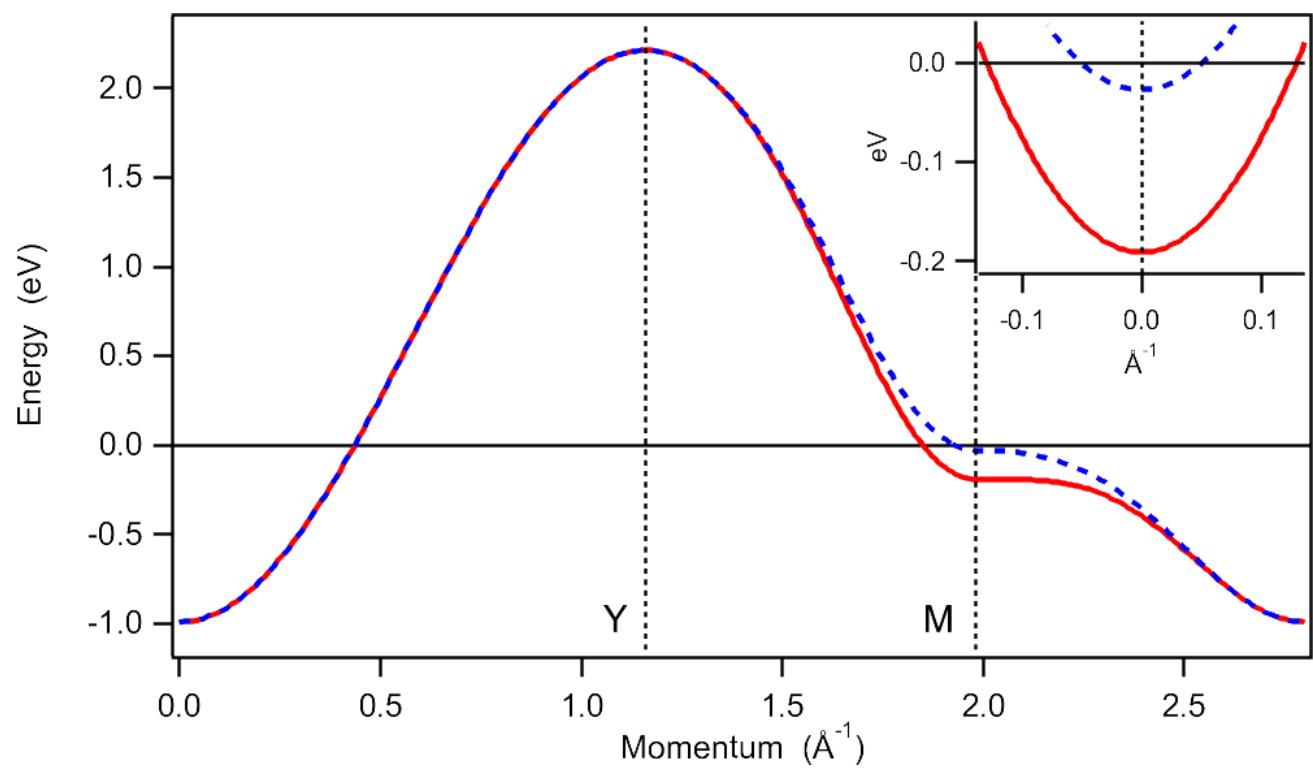
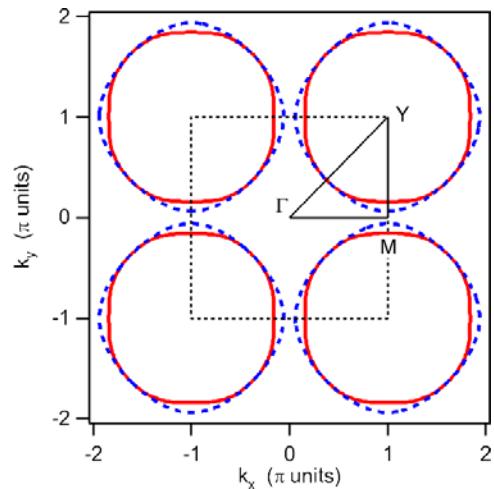
UD
 $< T_c$

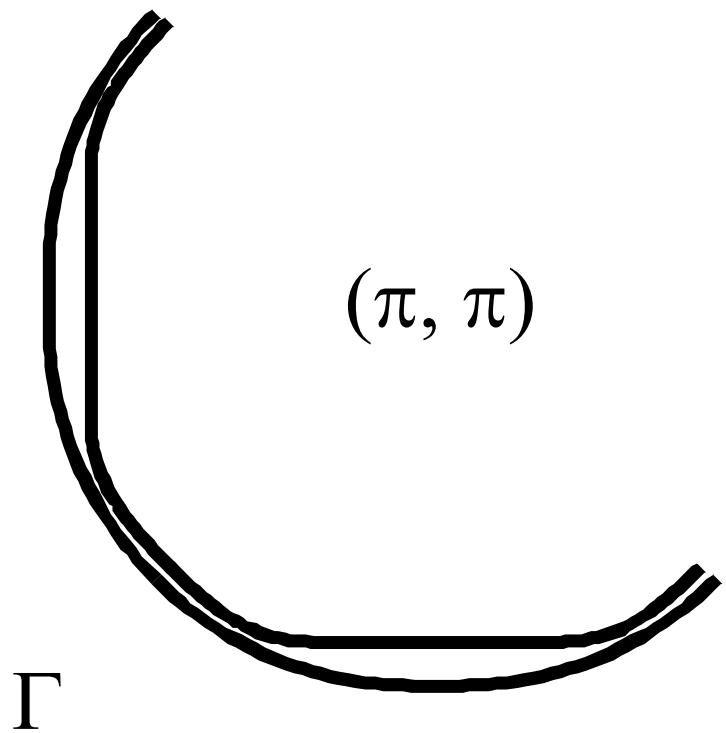
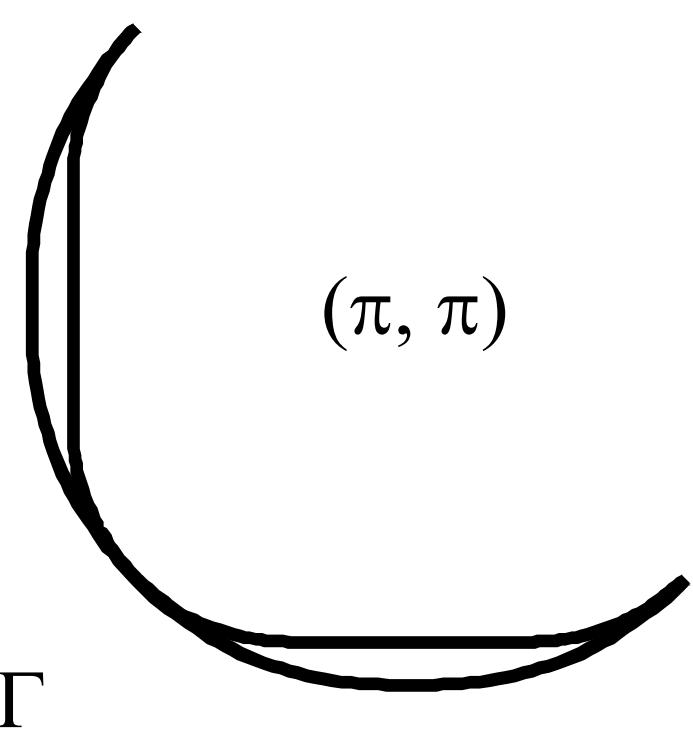


Bilayer splitting in OP Bi-2212 in normal state

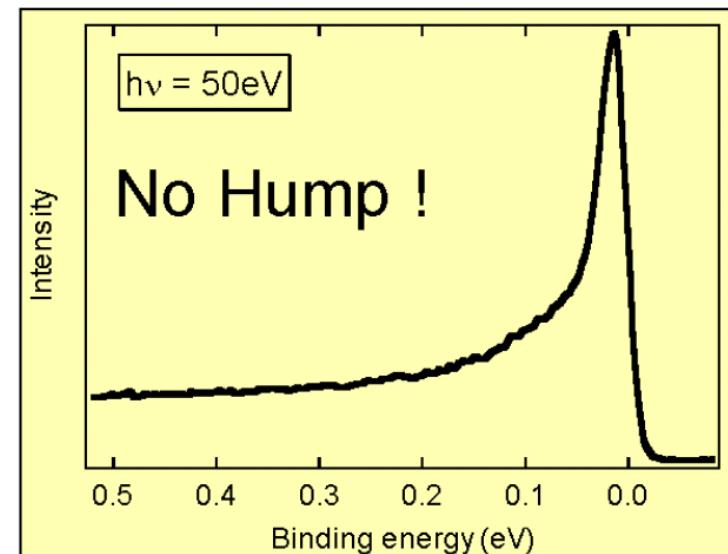
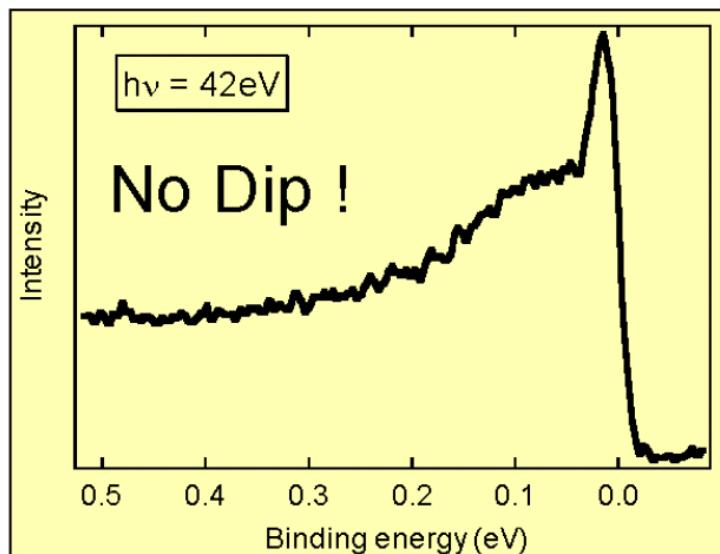
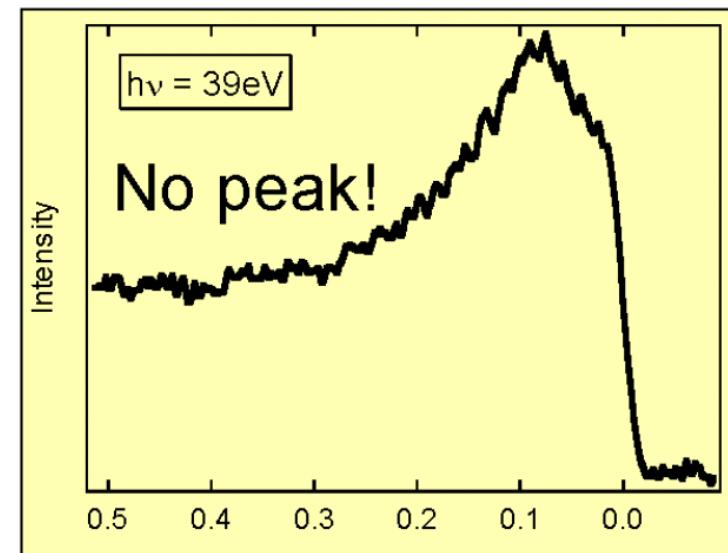
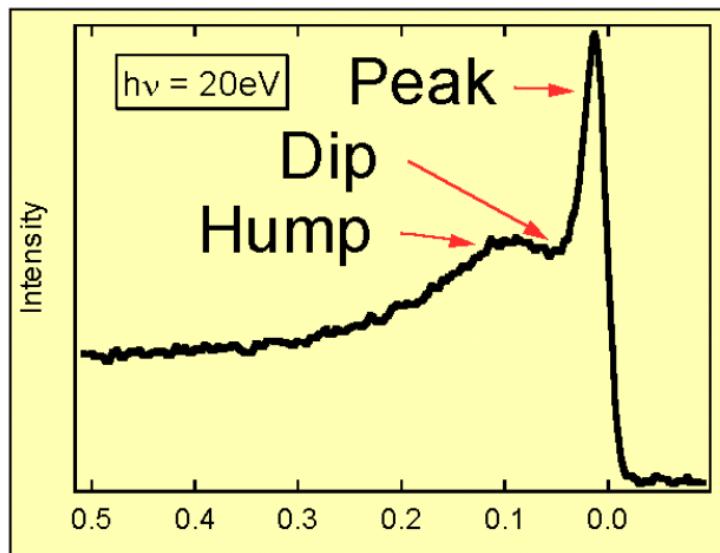


Bare band structure

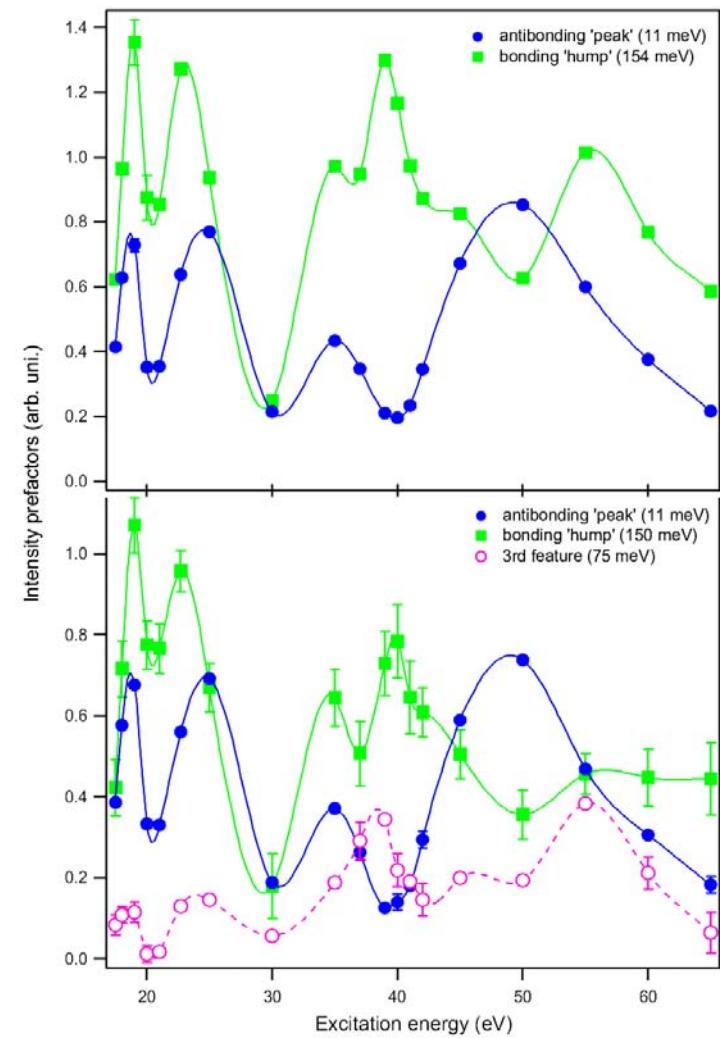
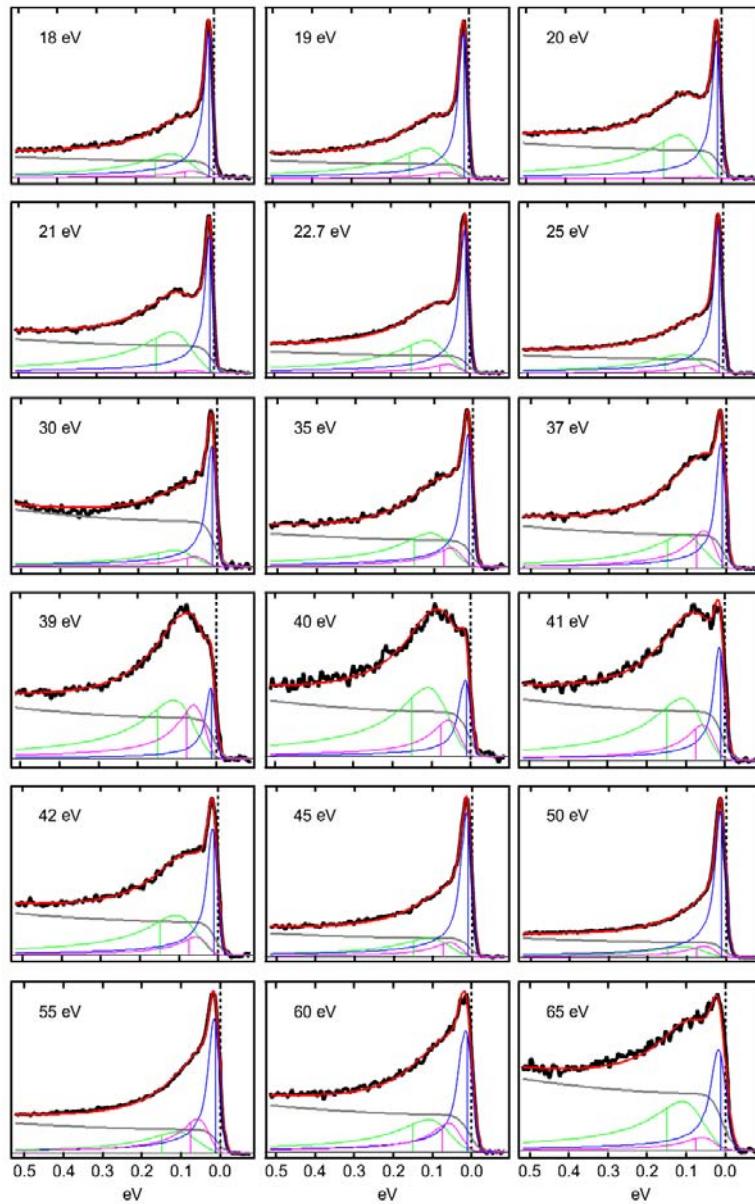




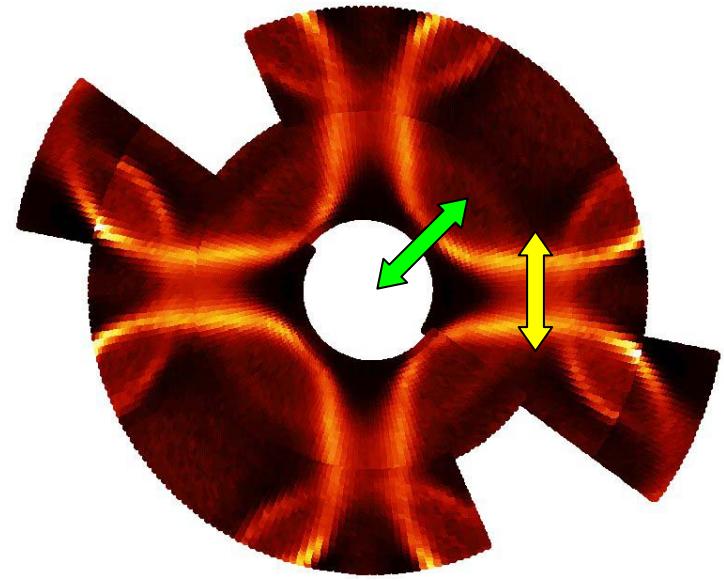
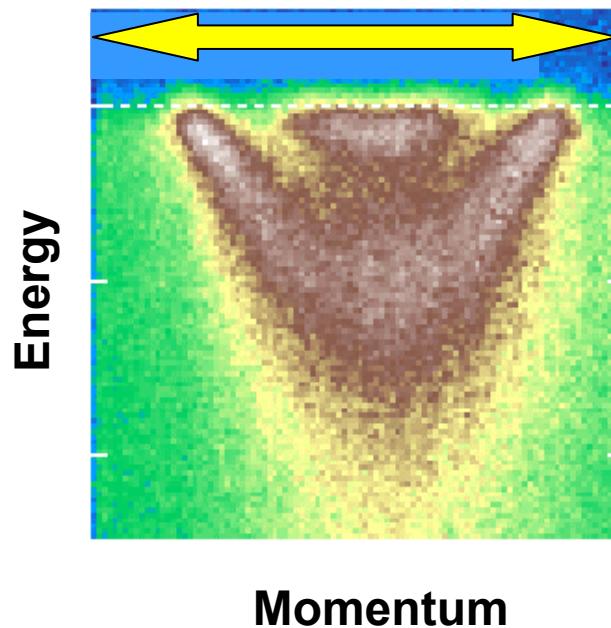
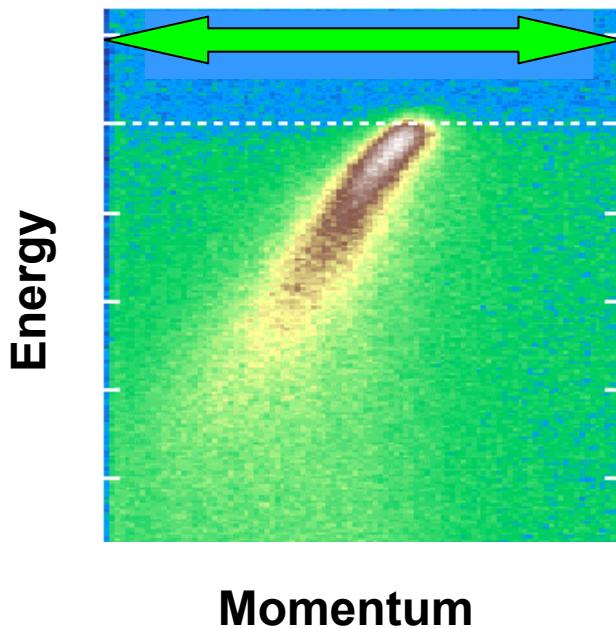
Excitation energy variation: PDH in OD



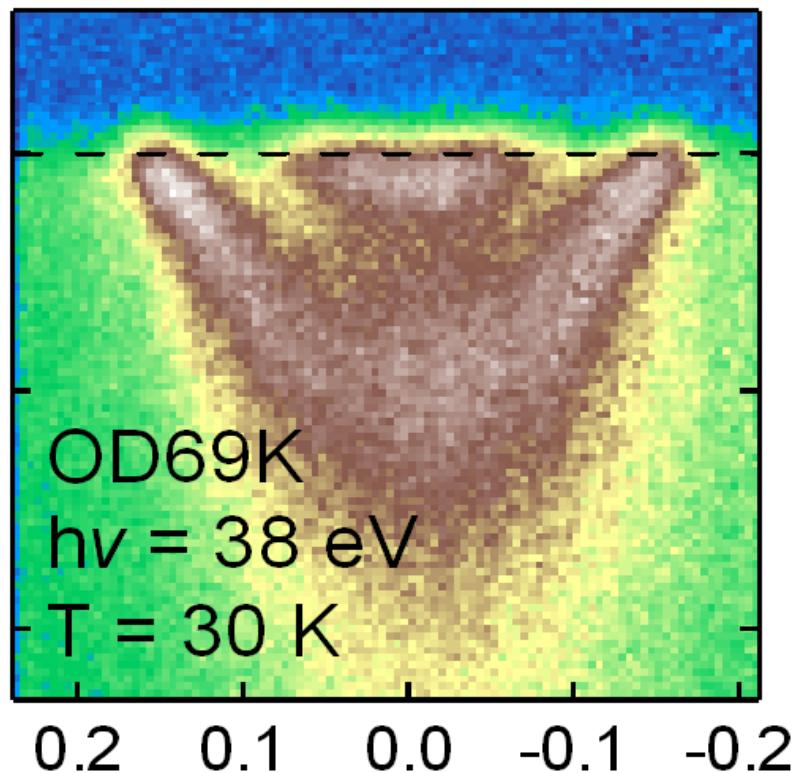
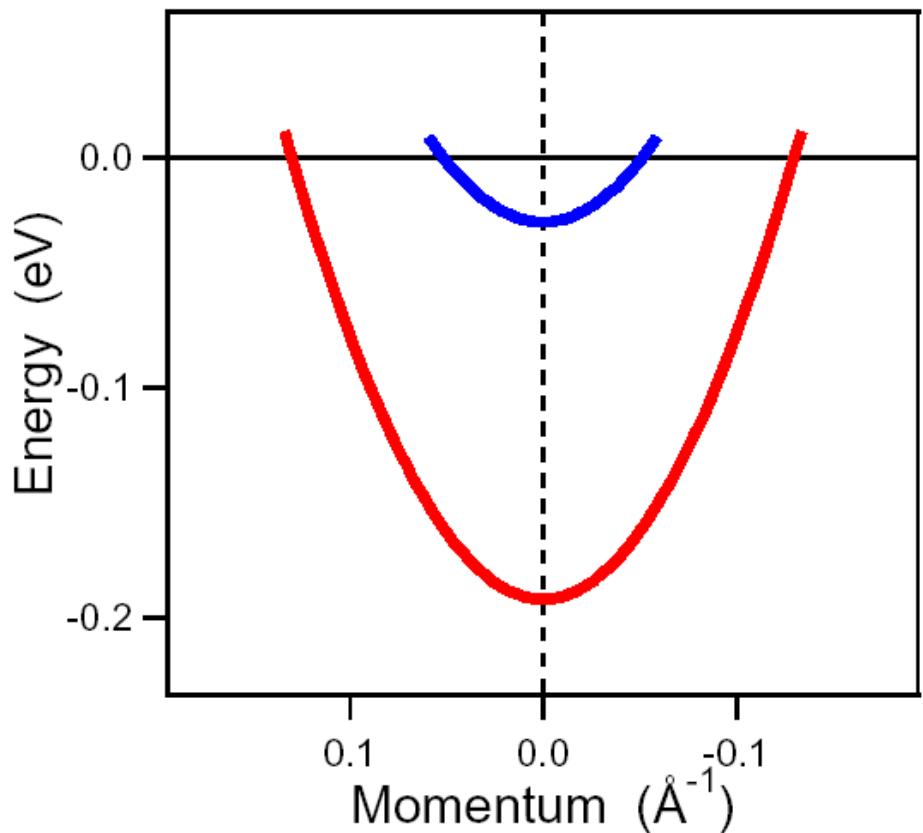
Excitation energy variation: PDH in OD



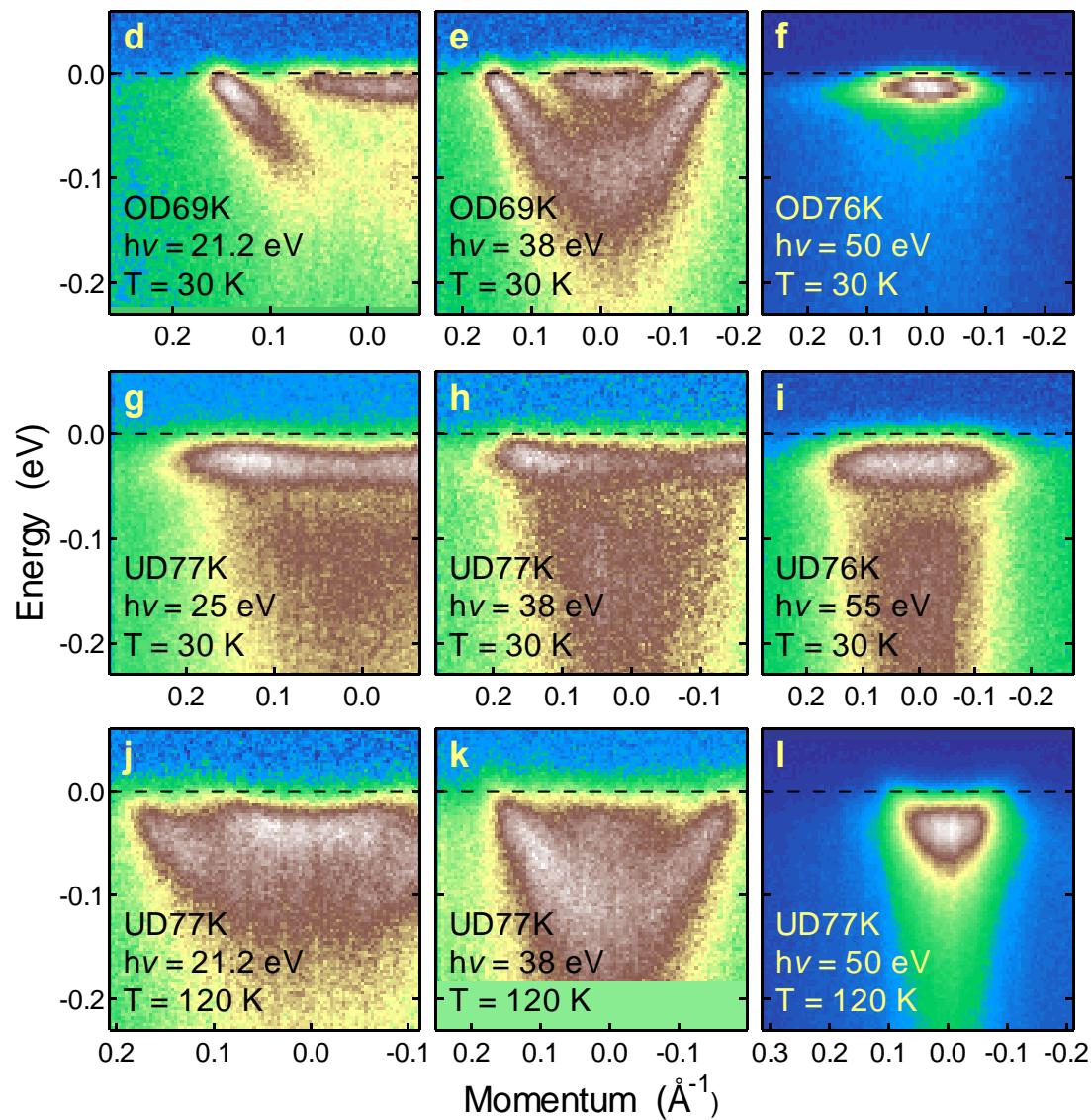
Key regions



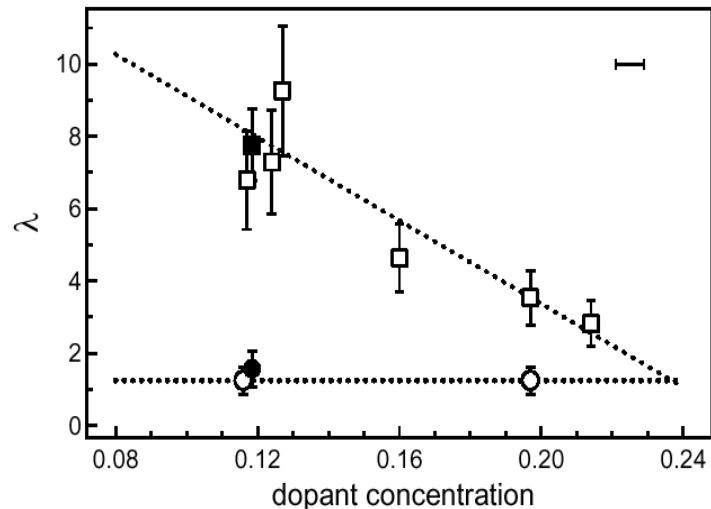
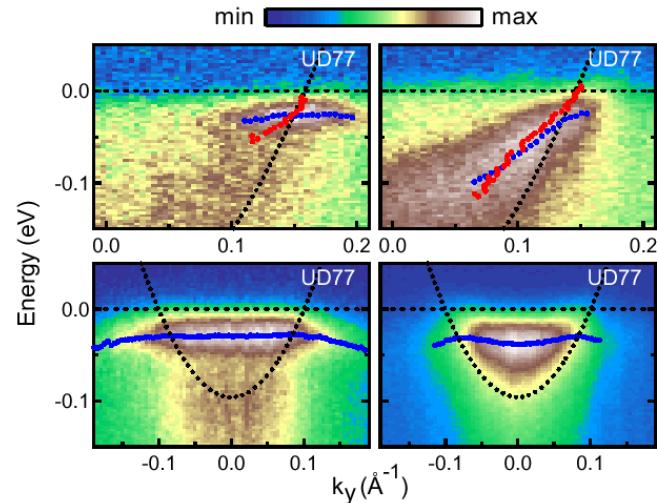
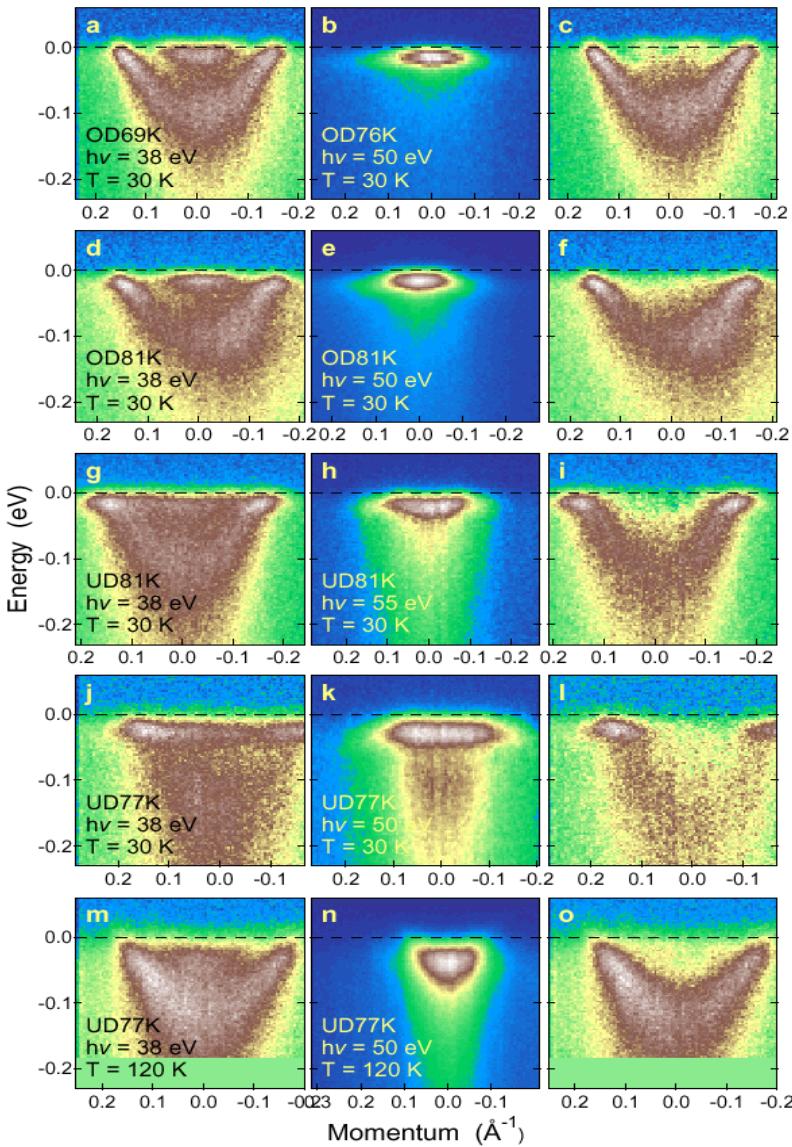
"XMY cut"



Interaction with a mode



Interaction with a mode



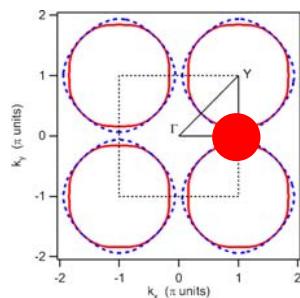
Electrons couple to ...

Doping dependence: UD \uparrow
OD \downarrow

Temperature dependence: $< T_c$

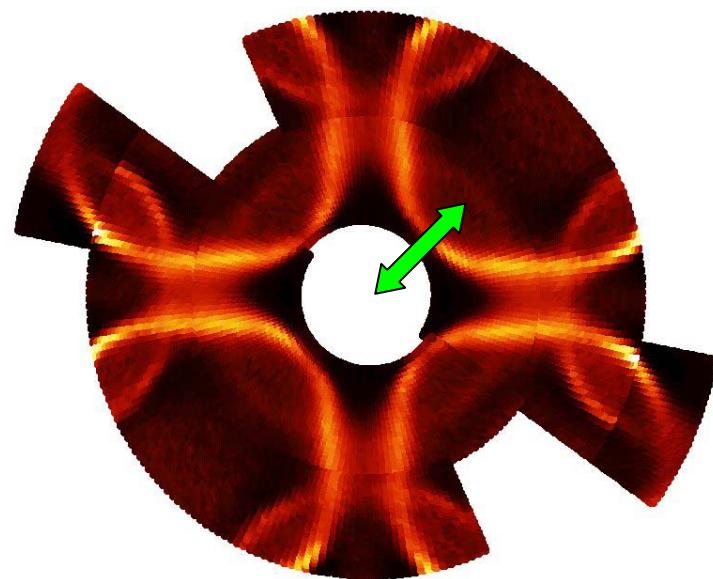
Energy ~ 40 meV

\mathbf{k} -dependence:

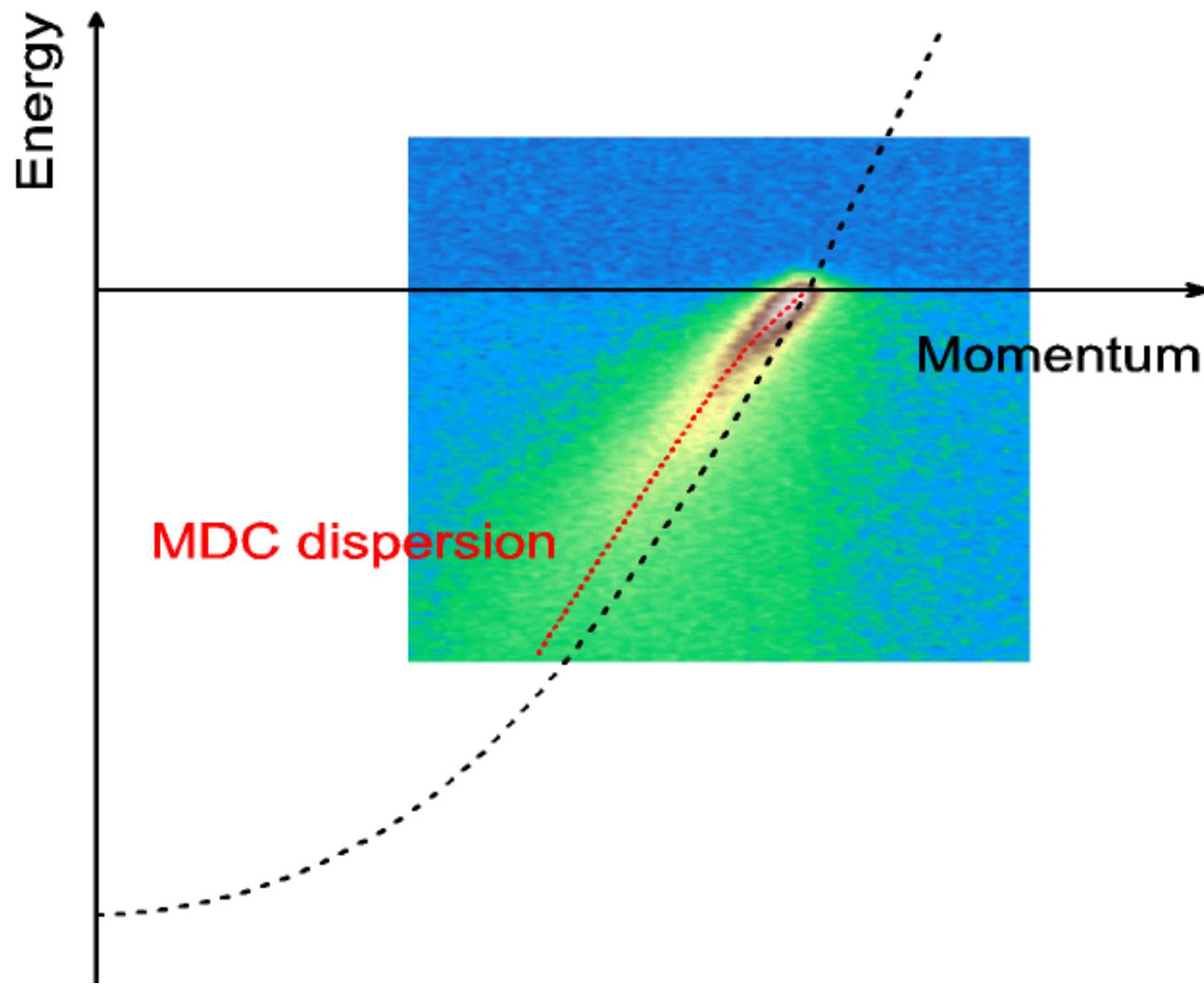


spin
fluctuations

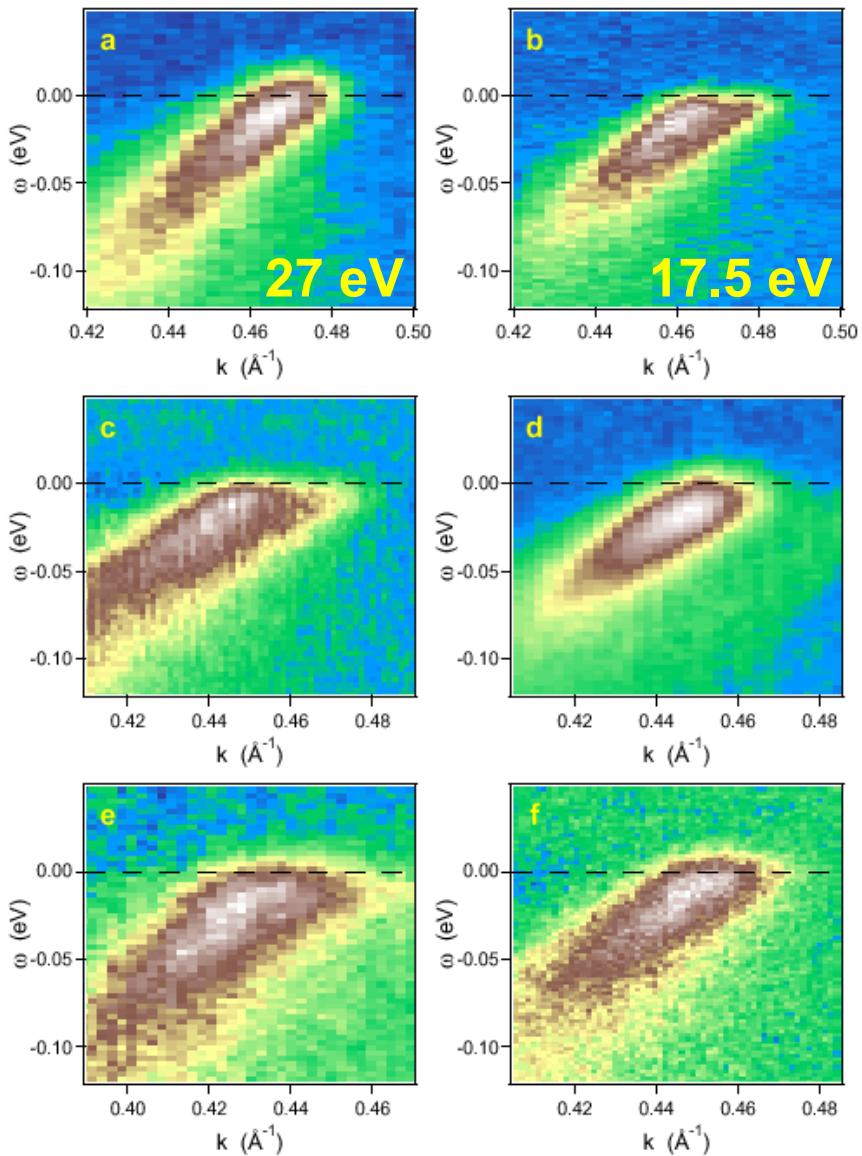
GX



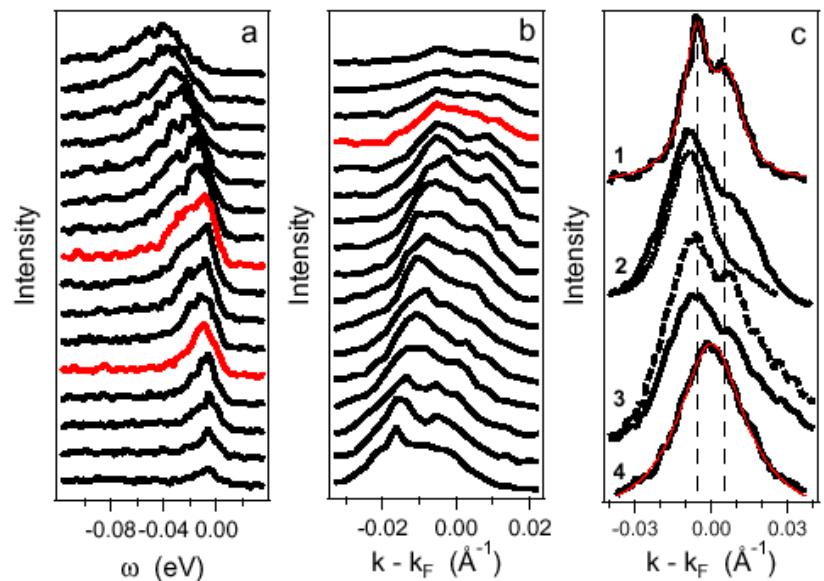
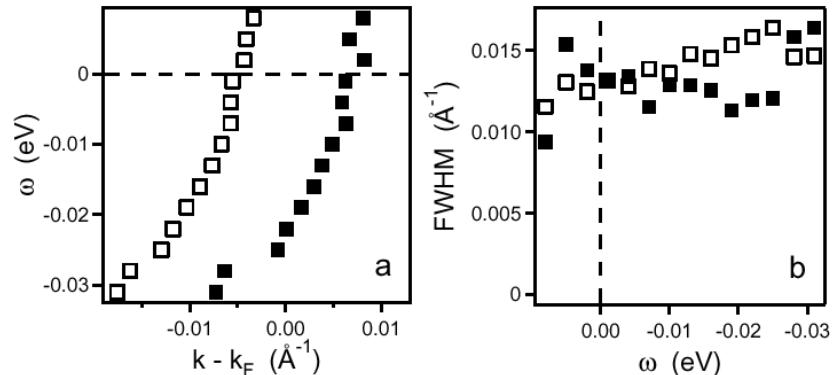
Basics: electron dispersion



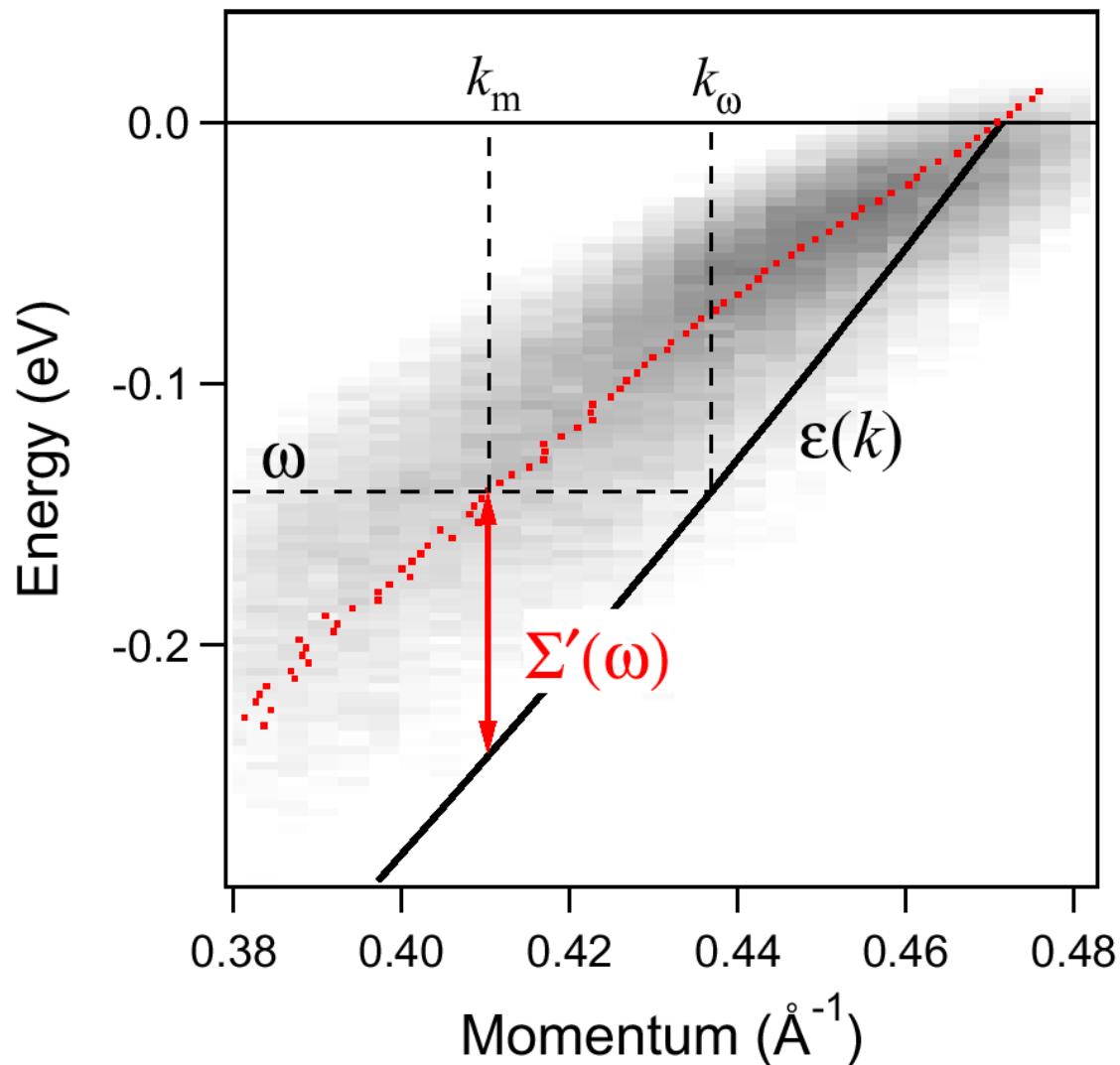
Nodal splitting



$$\Delta k = 0.012 \text{ } 1/\text{\AA}$$
$$\Delta \varepsilon = 50 \text{ meV}$$

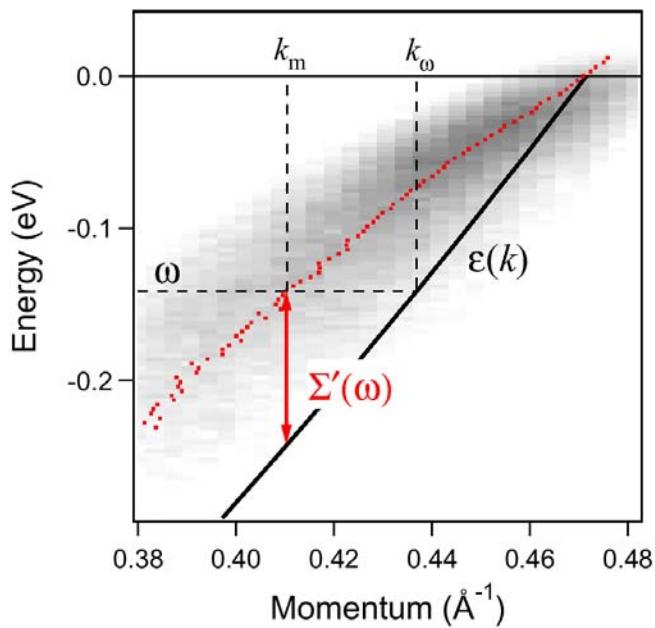


Bare dispersion



Self-energy approach

$$A(\omega, \mathbf{k}) = -\frac{1}{\pi} \frac{\Sigma''(\omega)}{(\omega - \varepsilon(\mathbf{k}) - \Sigma'(\omega))^2 + \Sigma''(\omega)^2}$$



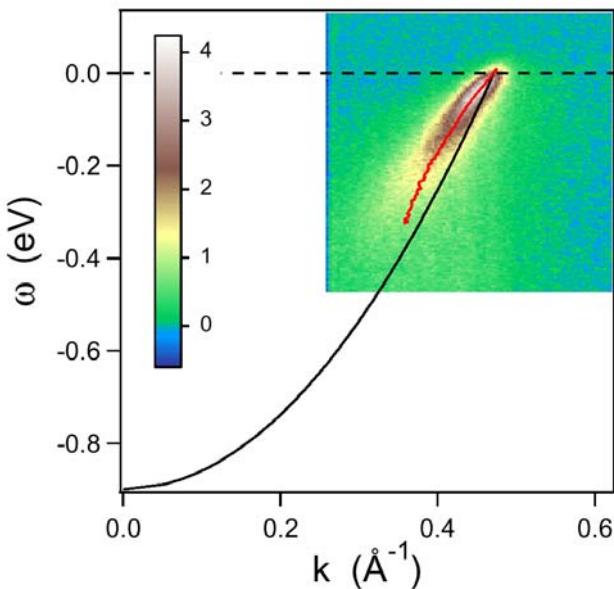
$$\Sigma'(\omega) = \omega - \varepsilon(k_m)$$

$$\Sigma''(\omega) = -v_F W(\omega)$$

Self-energy approach

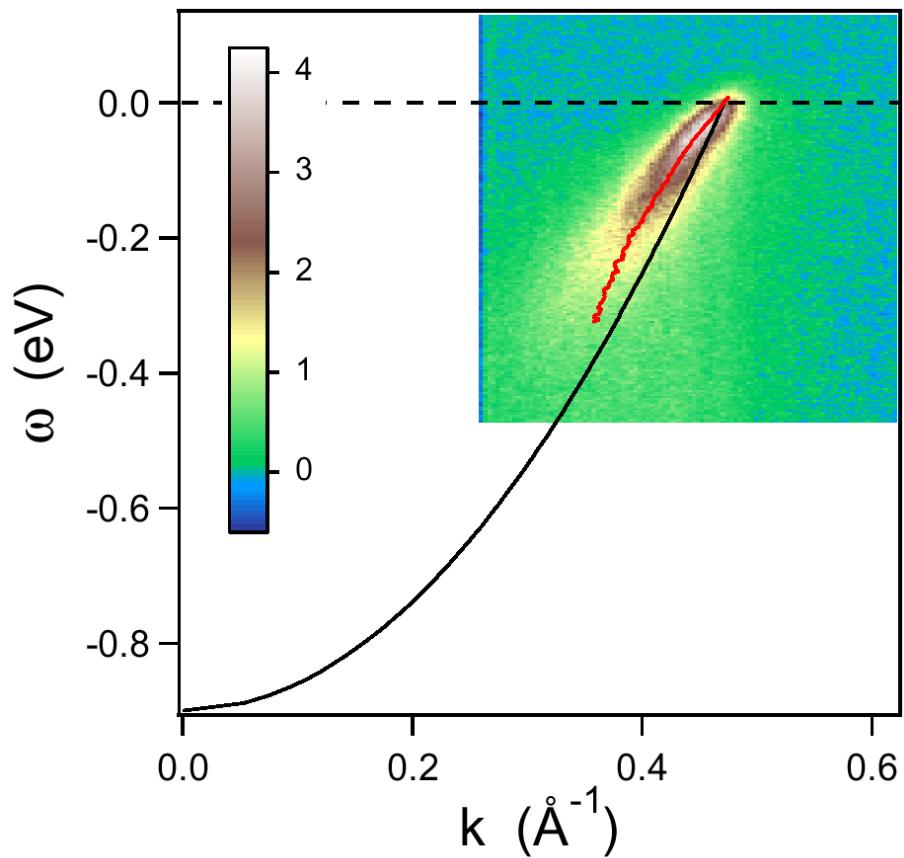
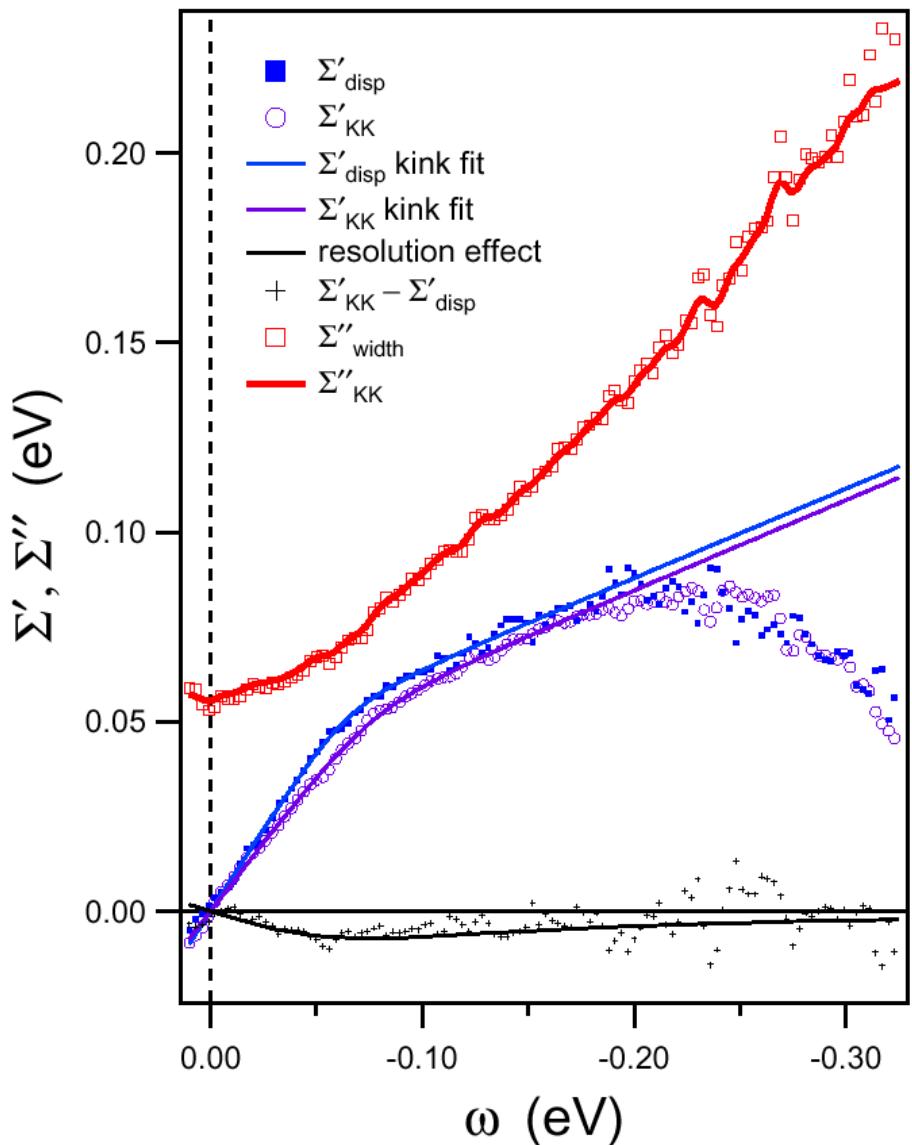
$$\Sigma'(\omega) = \frac{v_F}{2} (k_m^2(\omega) - k_F^2) + \omega,$$

$$\Sigma''(\omega) = -v_F W(\omega) \sqrt{k_m^2(\omega) - W^2(\omega)}.$$



$$\Sigma'(\omega) = \text{KK } \Sigma''(\omega)$$

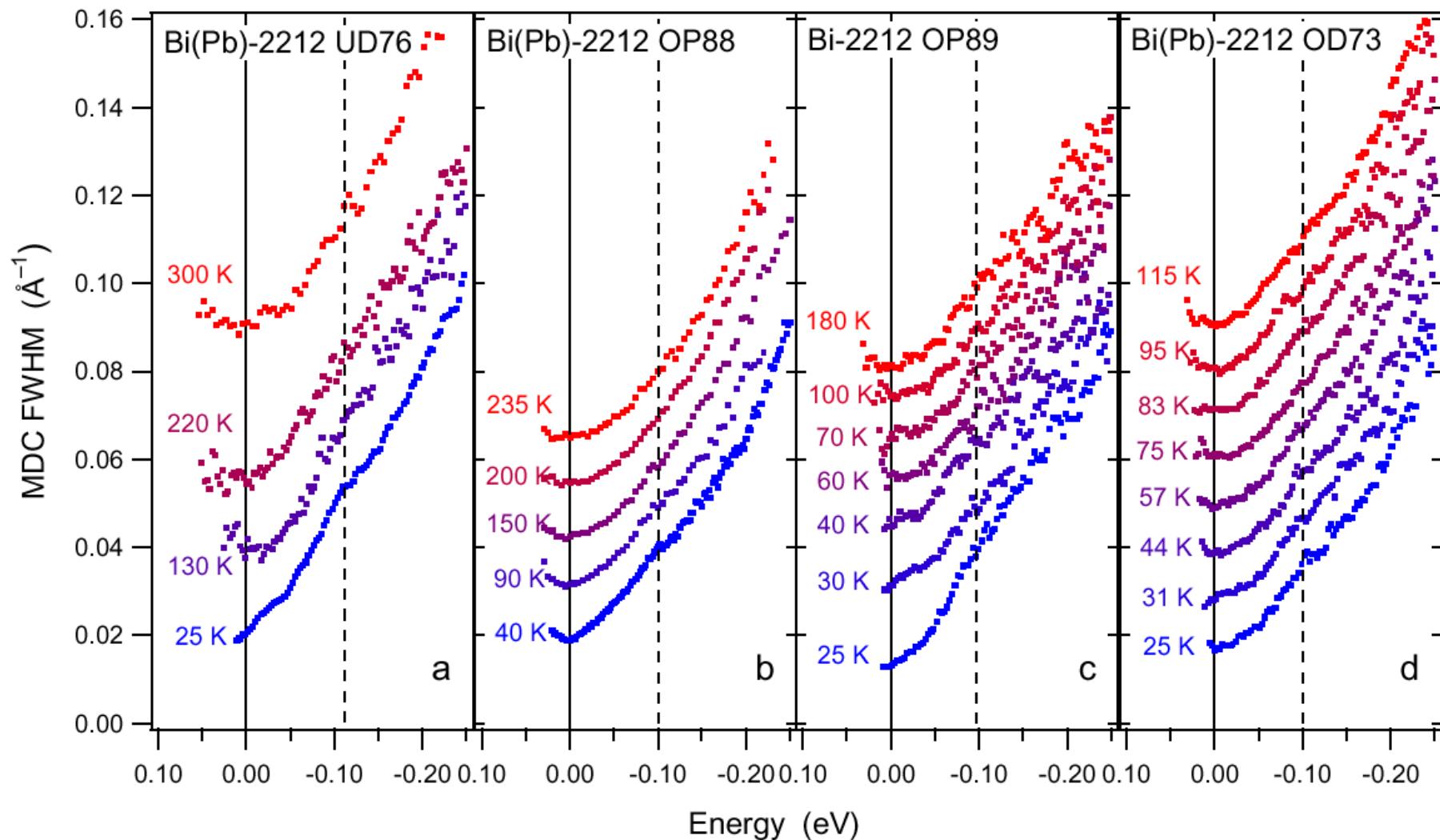
Bare dispersion



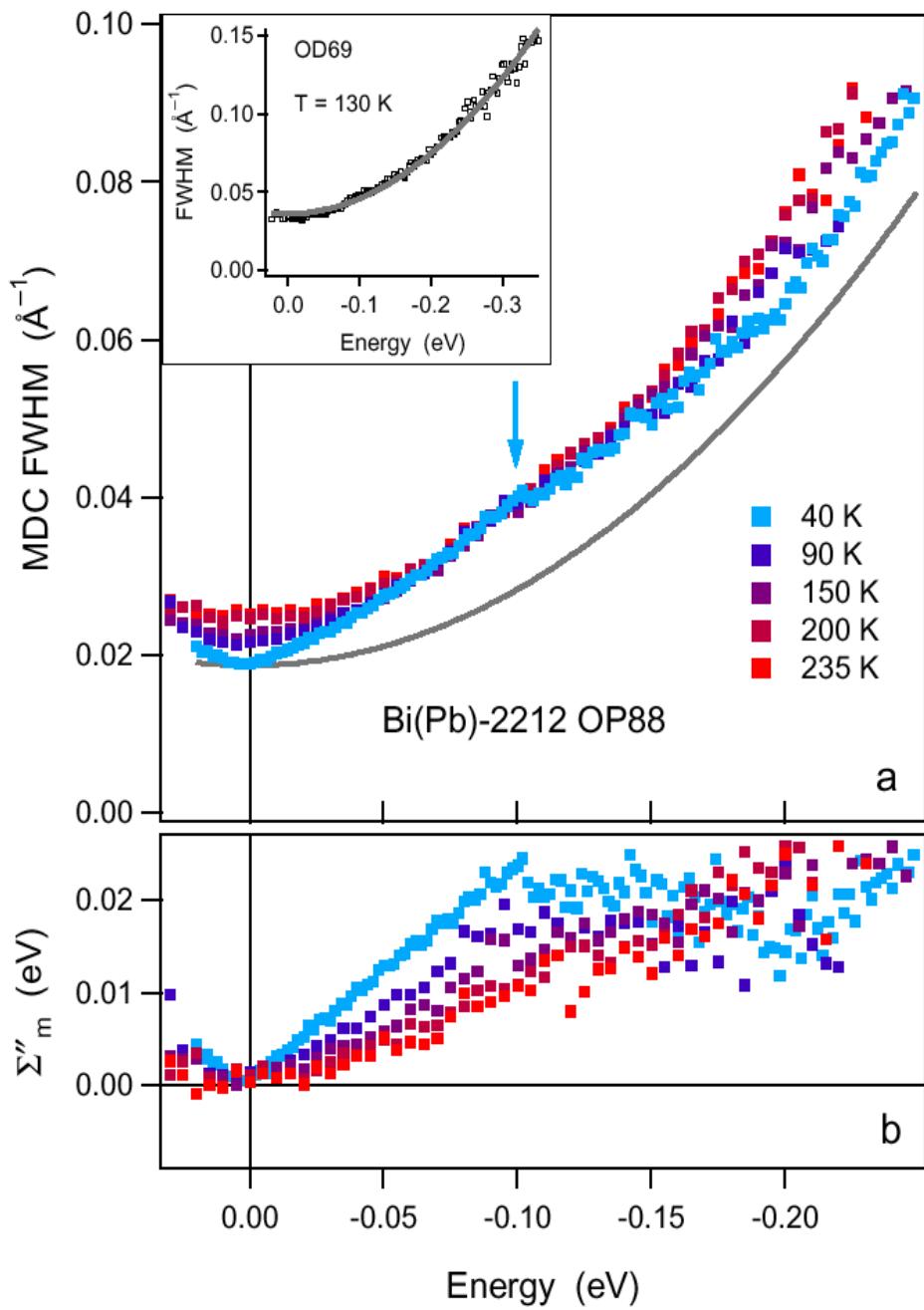
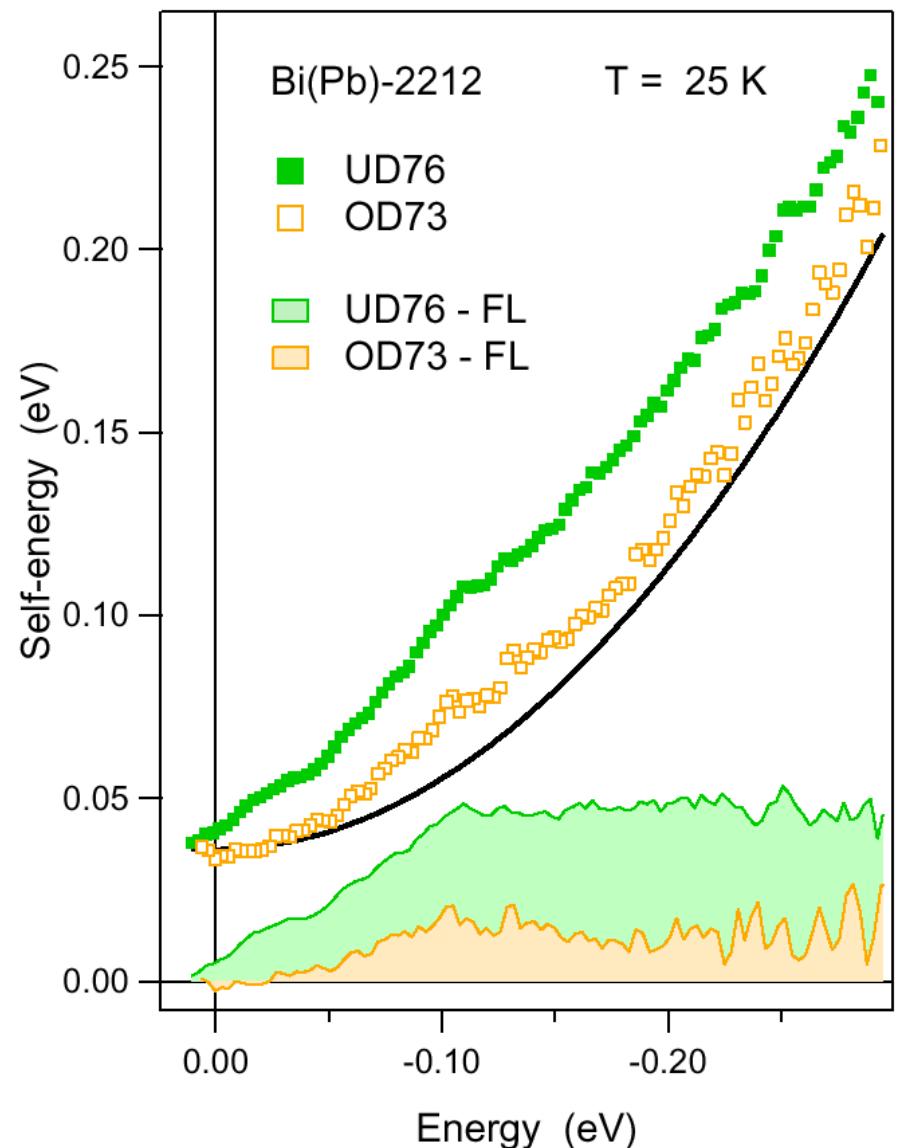
$$v_F = 3.82 \pm 0.17 \text{ eV\AA}$$

$$\lambda = 0.87 \pm 0.12$$

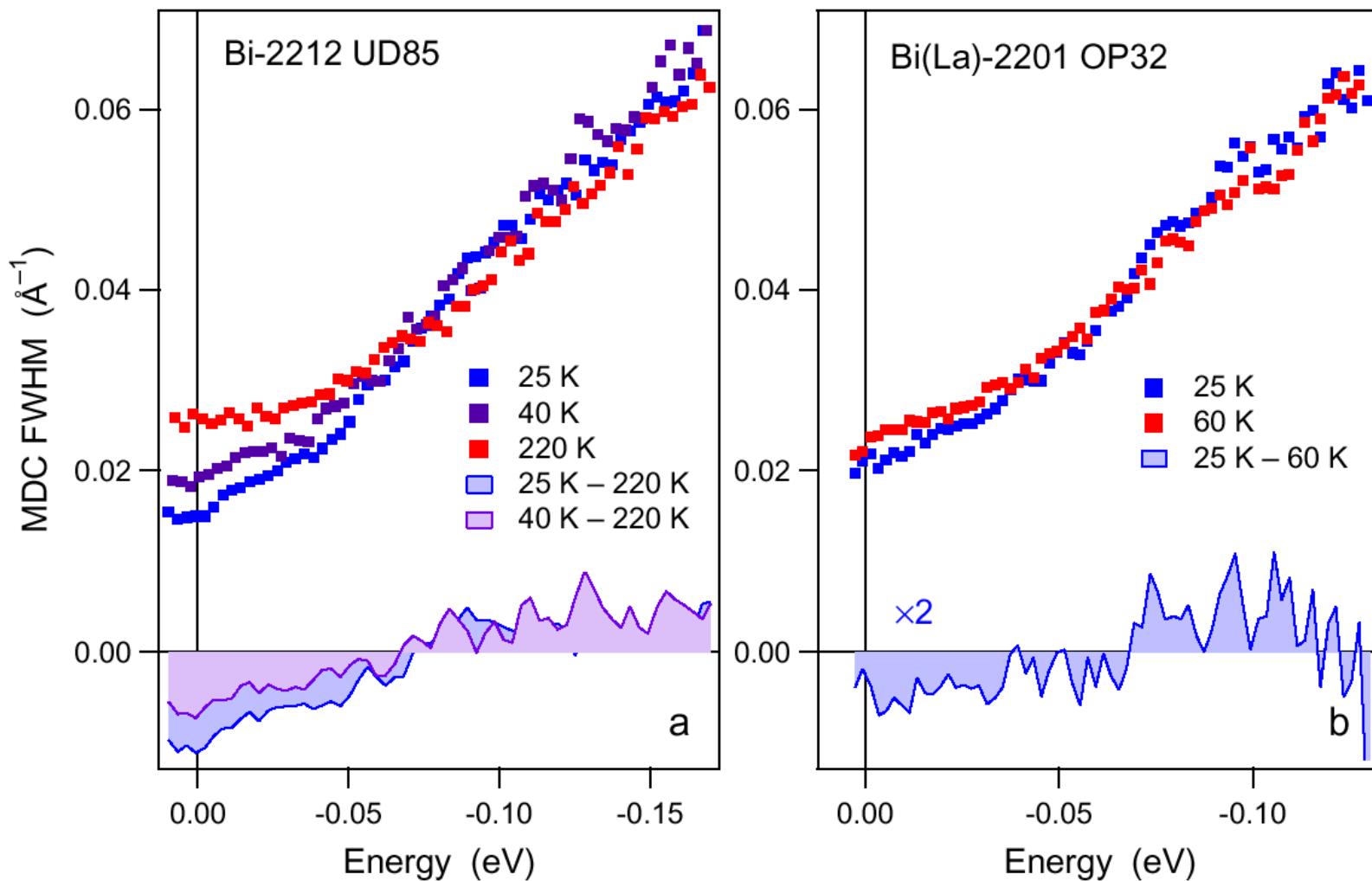
Scattering rate kink



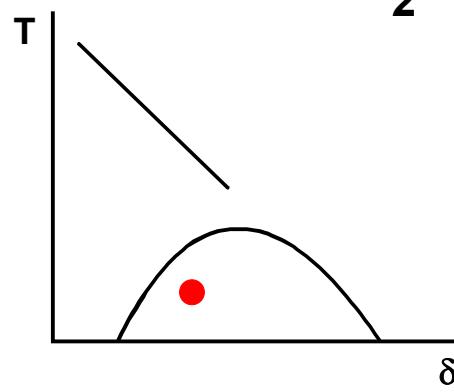
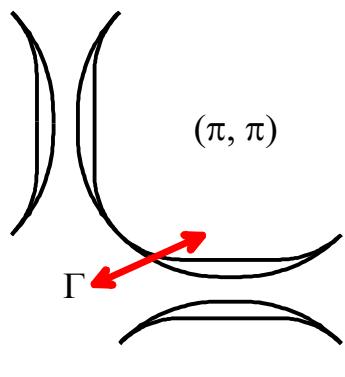
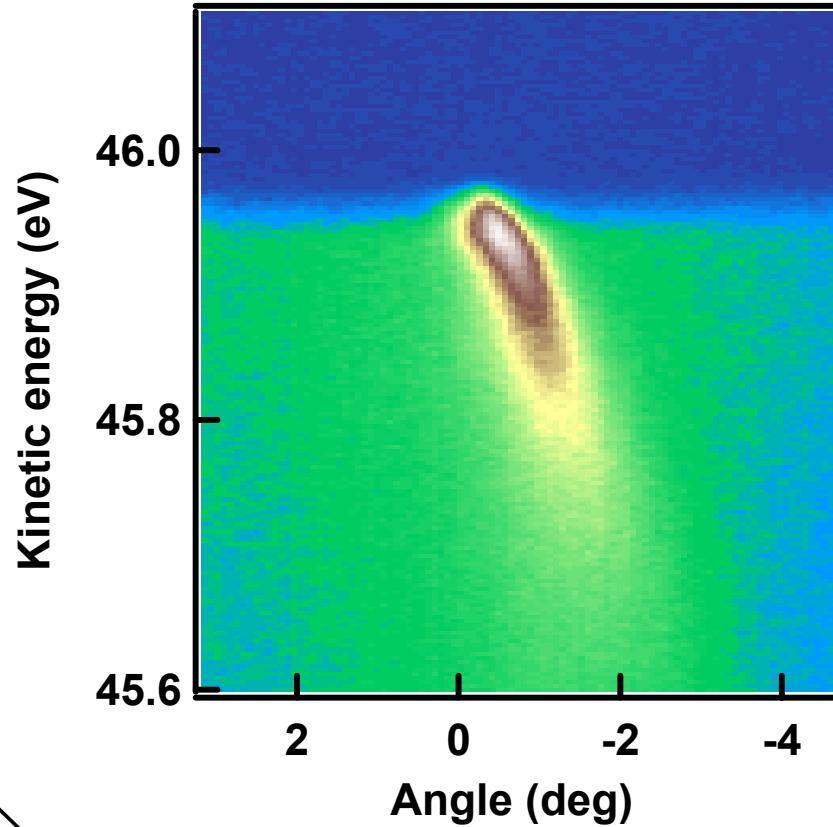
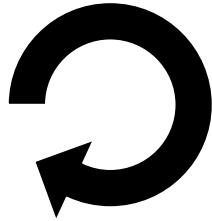
Scattering rate



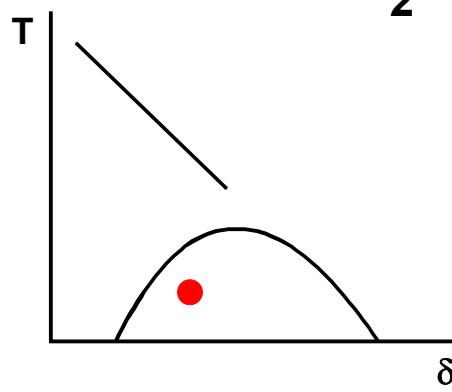
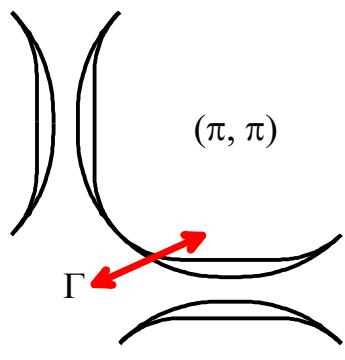
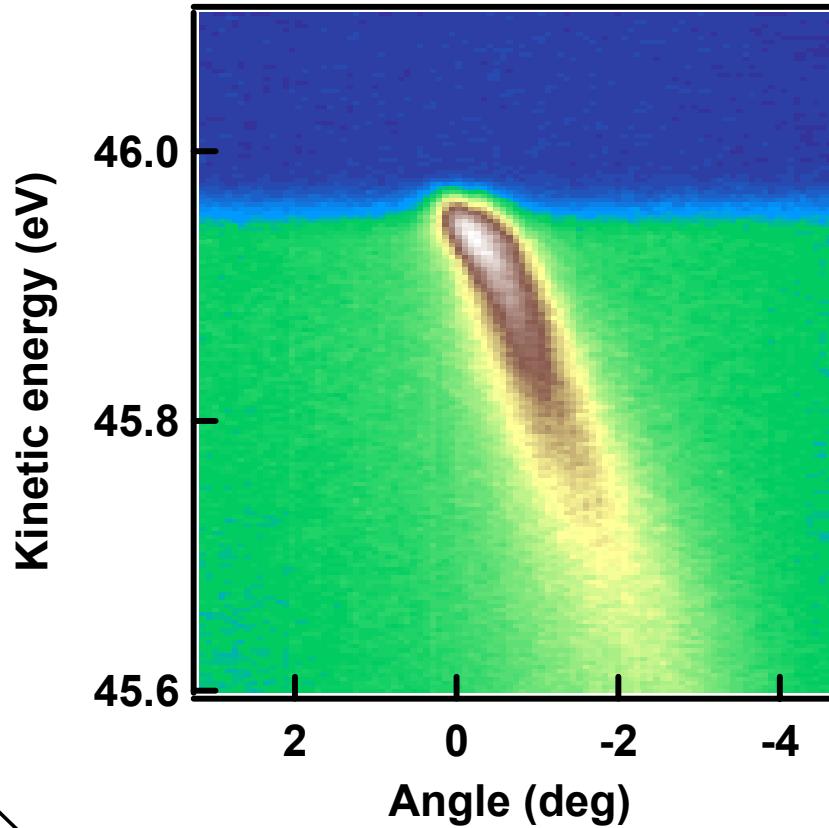
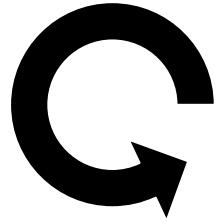
Scattering rate kink



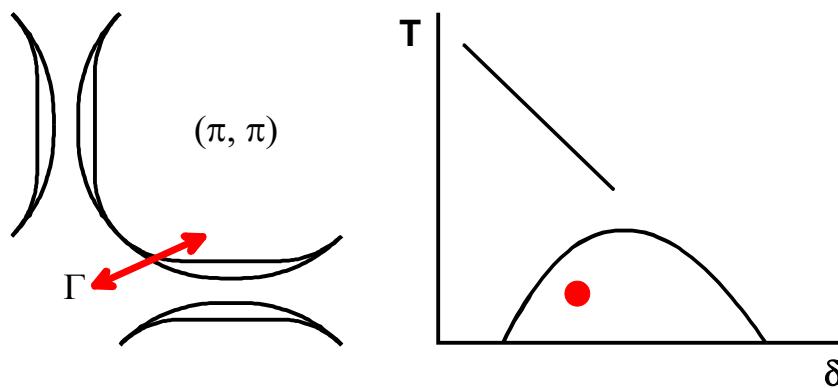
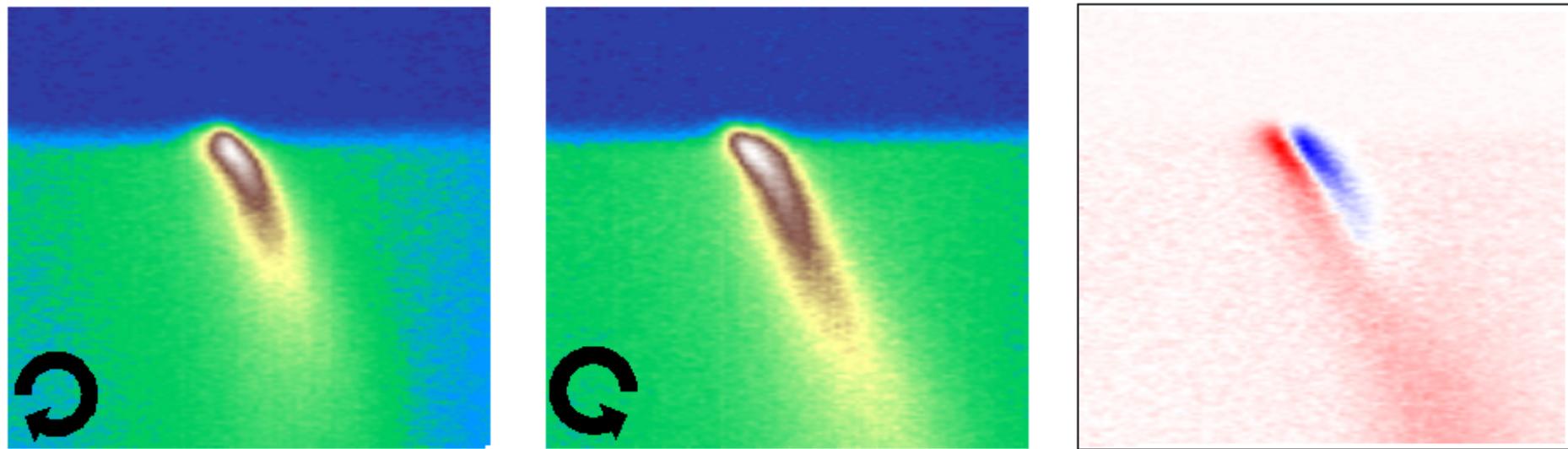
Circular dichroism in nodal region



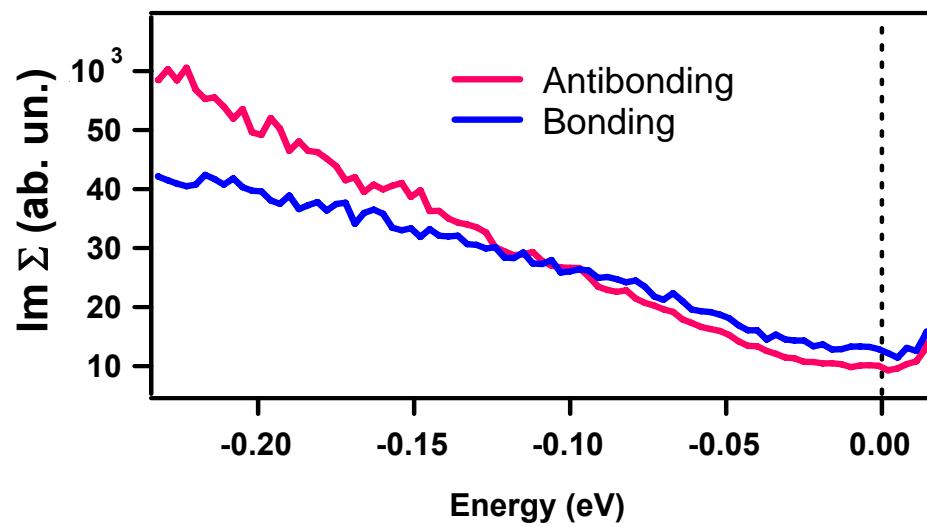
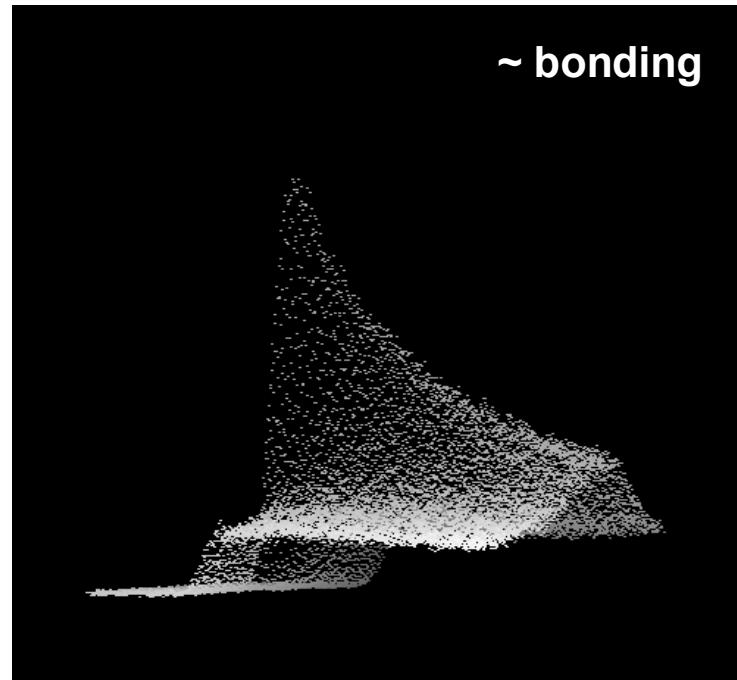
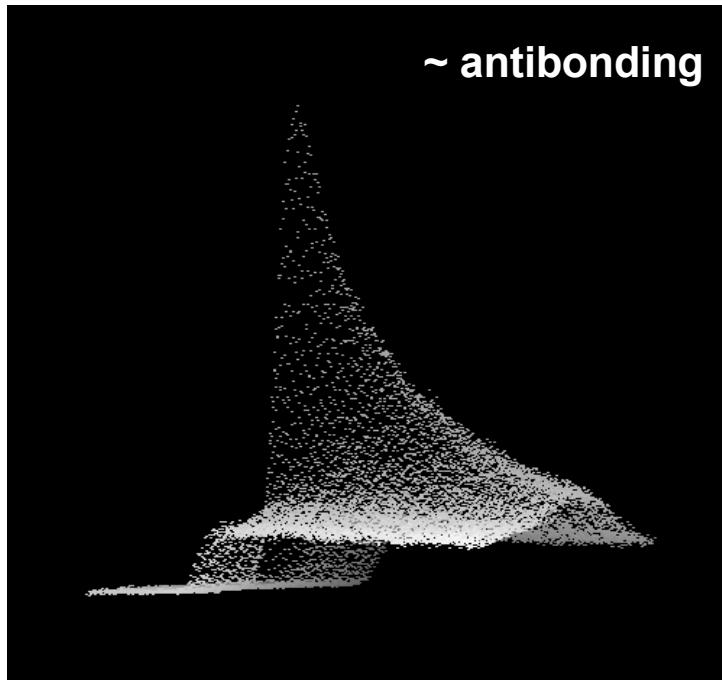
Circular dichroism in nodal region



Circular dichroism in nodal region



Odd scattering



Conclusions

- The spectral function analysis is applicable to the ARPES spectra from HTSC cuprates.
- Along the nodal direction **well defined quasiparticles** exist even for the underdoped Bi-2212 in the pseudogap state.
- **Two channels** in the scattering rate can be distinguished.
- The main doping independent contribution to the scattering can be well understood in terms of the conventional **Fermi liquid** model...
- ...while the additional doping dependent contribution has a **magnetic origin**.
- The magnetic contribution essentially increases with underdoping becoming dominant for the rest of the Brillouin zone and therefore determines the unusual properties of the cuprates in the superconducting and pseudo-gap phases.